

## Rigid Body Simulation

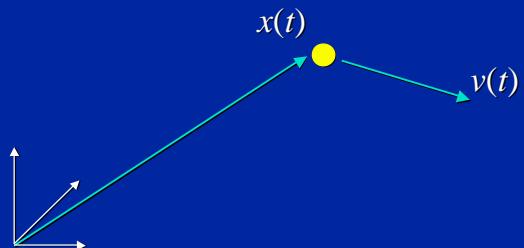
### Particle State

$$\mathbf{Y} = \begin{pmatrix} x(t) \\ v(t) \end{pmatrix}$$

$x(t)$        $v(t)$

$$\mathbf{Y} = \boxed{\phantom{0}} \boxed{\phantom{0}} \boxed{\phantom{0}} \boxed{\phantom{0}} \boxed{\phantom{0}} \boxed{\phantom{0}}$$

### Particle Motion



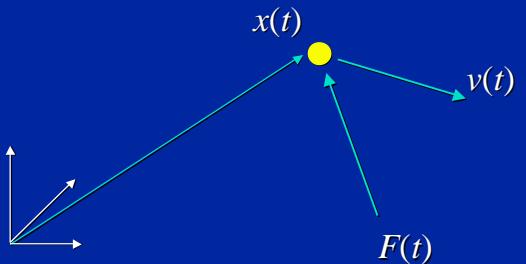
### State Derivative

$$\frac{d}{dt} \mathbf{Y} = \frac{d}{dt} \begin{pmatrix} x(t) \\ v(t) \end{pmatrix} = \begin{pmatrix} v(t) \\ F(t)/m \end{pmatrix}$$

$v(t)$        $F(t)/m$

$$\frac{d}{dt} \mathbf{Y} = \boxed{\phantom{0}} \boxed{\phantom{0}} \boxed{\phantom{0}} \boxed{\phantom{0}} \boxed{\phantom{0}} \boxed{\phantom{0}}$$

### Particle Dynamics

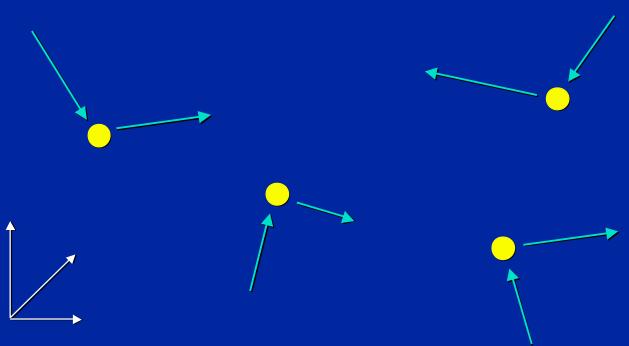


### State Derivative

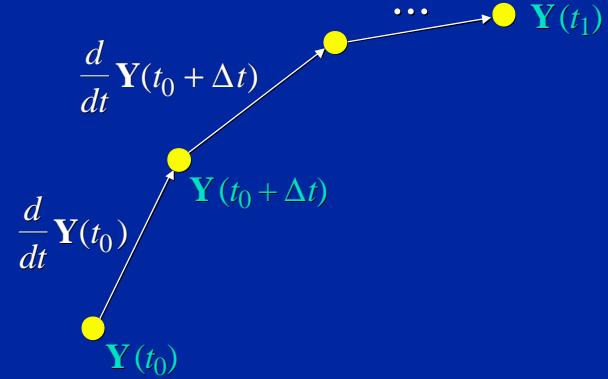
$$\frac{d}{dt} \mathbf{Y} = \frac{d}{dt} \begin{pmatrix} x_1(t) \\ v_1(t) \\ x_n(t) \\ v_n(t) \end{pmatrix} = \begin{pmatrix} v_1(t) \\ F_1(t)/m_1 \\ v_n(t) \\ F_n(t)/m_n \end{pmatrix}$$

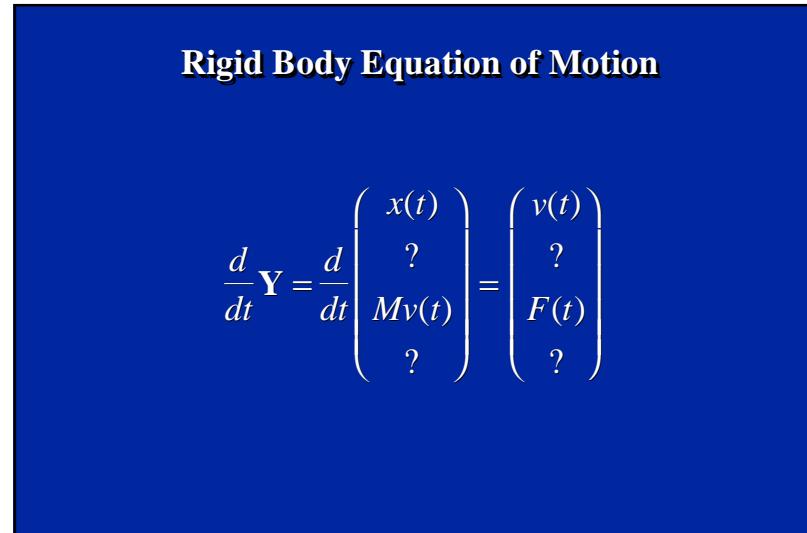
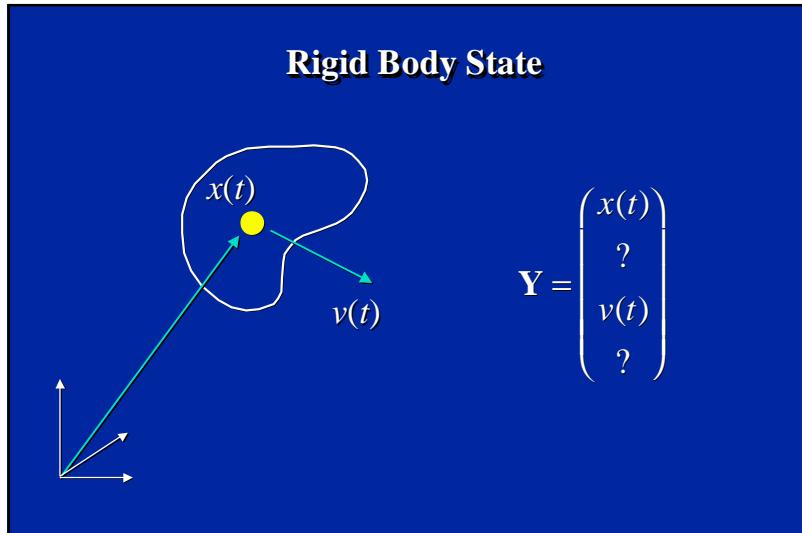
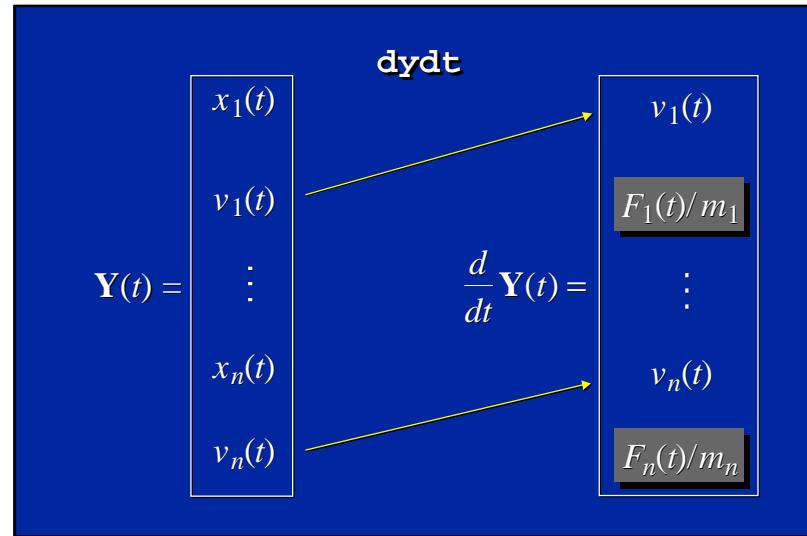
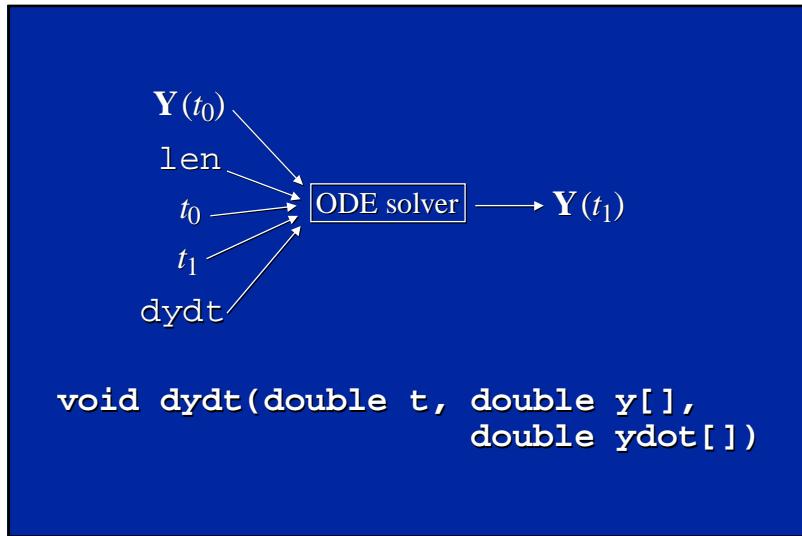
$$\frac{d}{dt} \mathbf{Y} = \boxed{\quad \quad} \dots 6n \text{ elements} \dots \boxed{\quad \quad}$$

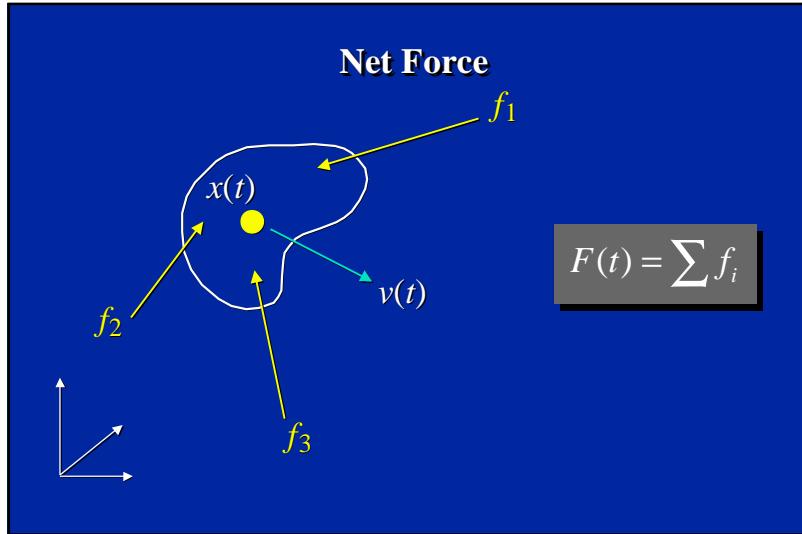
### Multiple Particles



### ODE solution







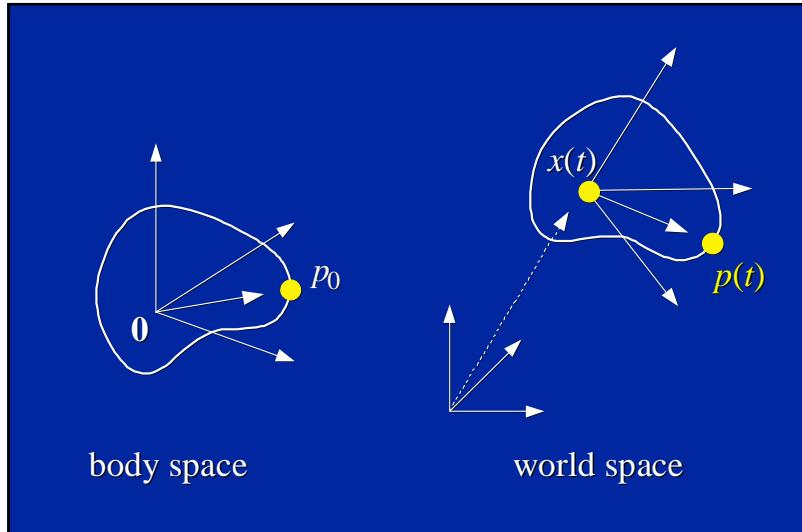
## Orientation

We represent orientation as a rotation matrix<sup>†</sup>  $R(t)$ . Points are transformed from body-space to world-space as:

$$p(t) = R(t)p_0 + x(t)$$


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<sup>†</sup>He's lying. Actually, we use quaternions.

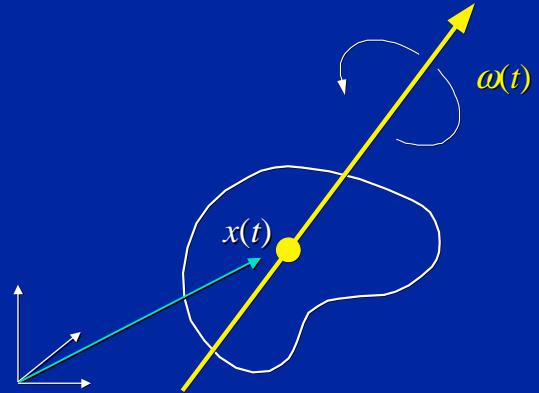


## Angular Velocity

We represent angular velocity as a vector  $\omega(t)$ , which encodes both the axis of the spin and the speed of the spin.

How are  $R(t)$  and  $\omega(t)$  related?

### Angular Velocity Definition



### Angular Velocity

$\dot{R}(t)$  and  $\omega(t)$  are related by

$$\frac{d}{dt} R(t) = \begin{pmatrix} 0 & -\omega_z(t) & \omega_y(t) \\ \omega_z(t) & 0 & -\omega_x(t) \\ -\omega_y(t) & \omega_x(t) & 0 \end{pmatrix} R(t)$$

( $\omega(t)^*$  is a shorthand for the above matrix)

### Rigid Body Equation of Motion

$$\frac{d}{dt} \mathbf{Y} = \frac{d}{dt} \begin{pmatrix} x(t) \\ R(t) \\ Mv(t) \\ \langle \omega(t) \rangle \end{pmatrix} = \begin{pmatrix} v(t) \\ \omega(t)^* R(t) \\ F(t) \\ ? \end{pmatrix}$$

Need to relate  $\dot{\omega}(t)$  and mass distribution to  $F(t)$ .

### Inertia Tensor

$$I(t) = \begin{pmatrix} I_{xx} & I_{xy} & I_{xz} \\ I_{yx} & I_{yy} & I_{yz} \\ I_{zx} & I_{zy} & I_{zz} \end{pmatrix}$$

diagonal terms

$$I_{xx} = M \int_V (y^2 + z^2) dV$$

$$I_{xy} = -M \int_V xy dV$$

off-diagonal terms

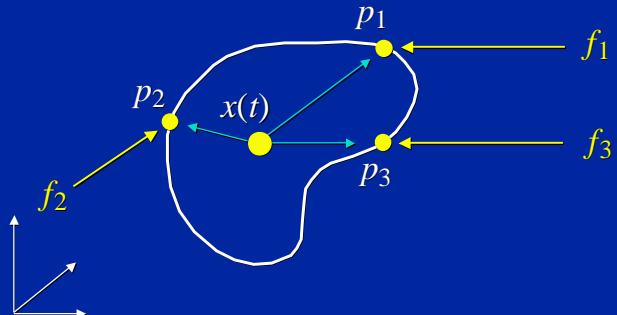
## Rigid Body Equation of Motion

$$\frac{d}{dt} \mathbf{Y} = \frac{d}{dt} \begin{pmatrix} x(t) \\ R(t) \\ \boxed{Mv(t)} \\ \boxed{I(t)\omega(t)} \end{pmatrix} = \begin{pmatrix} v(t) \\ \omega(t)^* R(t) \\ F(t) \\ \tau(t) \end{pmatrix}$$

$P(t)$  – linear momentum

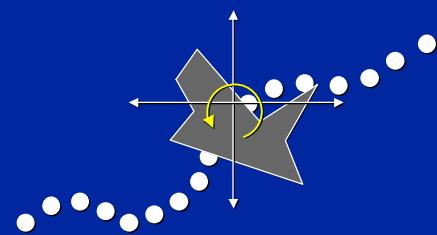
$L(t)$  – angular momentum

## Net Torque



$$\tau(t) = \sum (p_i - x(t)) \times f_i$$

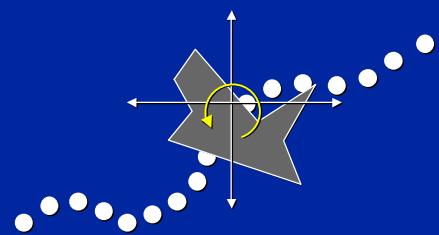
Inertia Tensors Vary in World Space...



$$I_{xx} = M \int_V (y^2 + z^2) dV$$

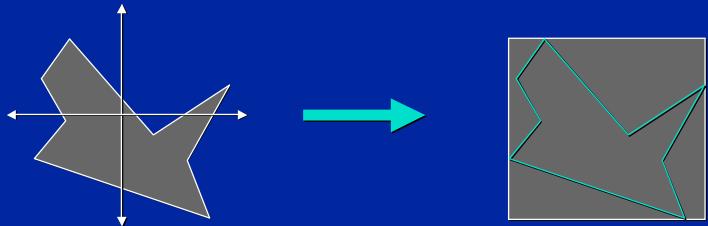
$$I_{xy} = -M \int_V xy dV$$

... but are Constant in Body Space



$$I(t) = R(t) I_{\text{body}} R(t)^T$$

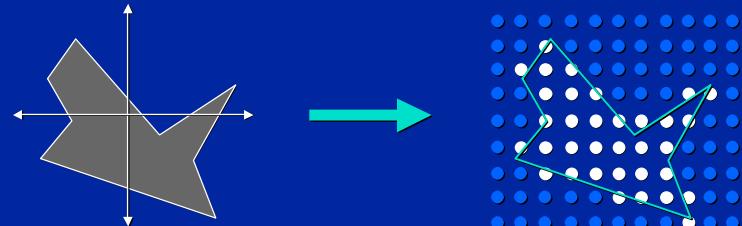
### Approximating $I_{\text{body}}$ —Bounding Boxes



Pros: Simple.

Cons: Bounding box may not be a good fit.  
Inaccurate.

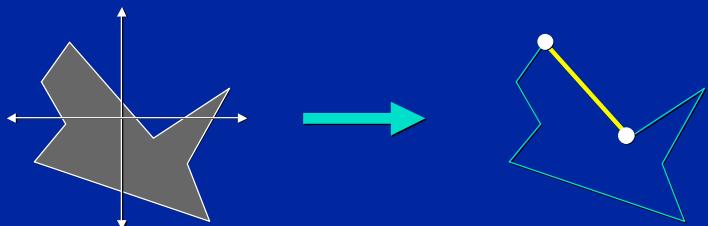
### Approximating $I_{\text{body}}$ —Point Sampling



Pros: Simple, fairly accurate, no B-rep needed.

Cons: Expensive, requires volume test.

### Computing $I_{\text{body}}$ —Green's Theorem (Twice!)



Pros: Simple, exact, no volumes needed.

Cons: Requires B-rep.

Code: <http://www.acm.org/jgt/papers/Mirtich96>

### Computing $I_{\text{body}}$ at the center of mass

The inertia tensor should be relative to the body's center of mass, e.g:

$$I_{xx} = I'_{xx} - m(r_y^2 + r_z^2)$$

$$I_{xy} = I'_{xy} - mr_x r_y$$

## Rigid Body Equation of Motion

$$\frac{d}{dt} \mathbf{Y} = \frac{d}{dt} \begin{pmatrix} x(t) \\ R(t) \\ \boxed{Mv(t)} \\ \boxed{I(t)\omega(t)} \end{pmatrix} = \begin{pmatrix} v(t) \\ \omega(t)^* R(t) \\ F(t) \\ \tau(t) \end{pmatrix}$$

$P(t)$  – linear momentum

$L(t)$  – angular momentum