

Computing in carbon

Basic elements of neuroelectronics

- membranes
- ion channels
- wiring

Elementary neuron models

- conductance based
- modelers' alternatives

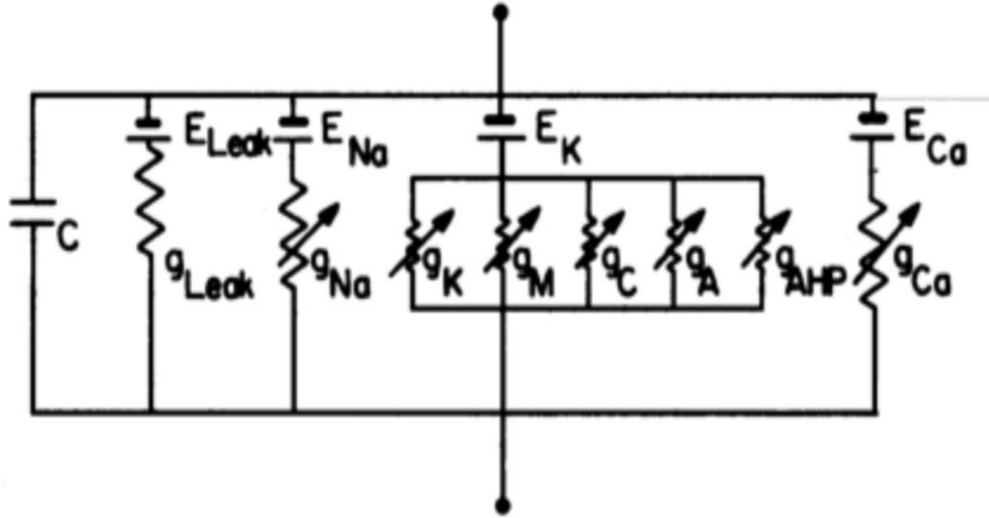
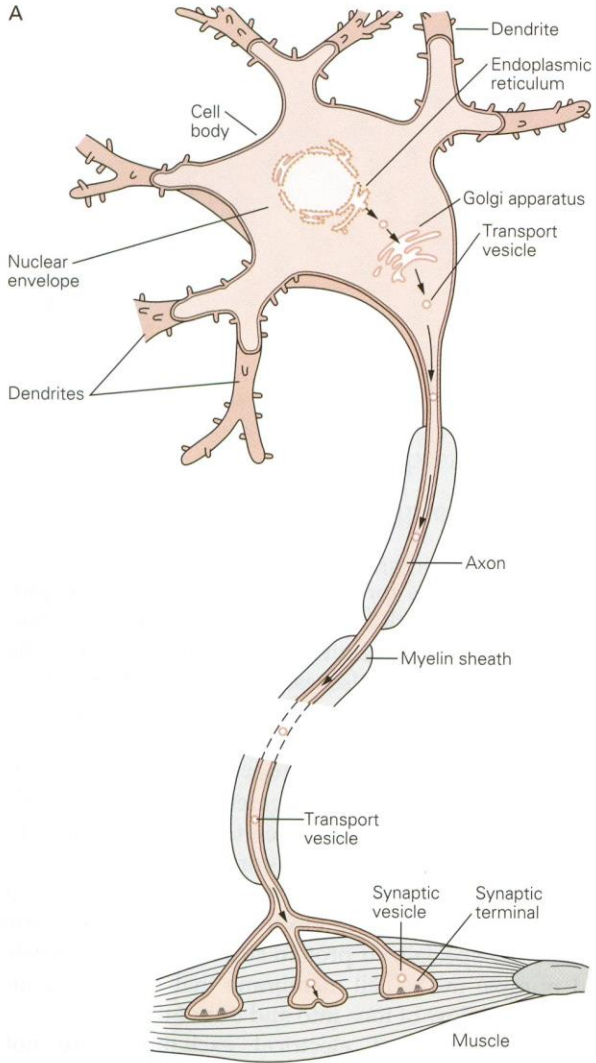
Wires

- signal propagation
- processing in dendrites

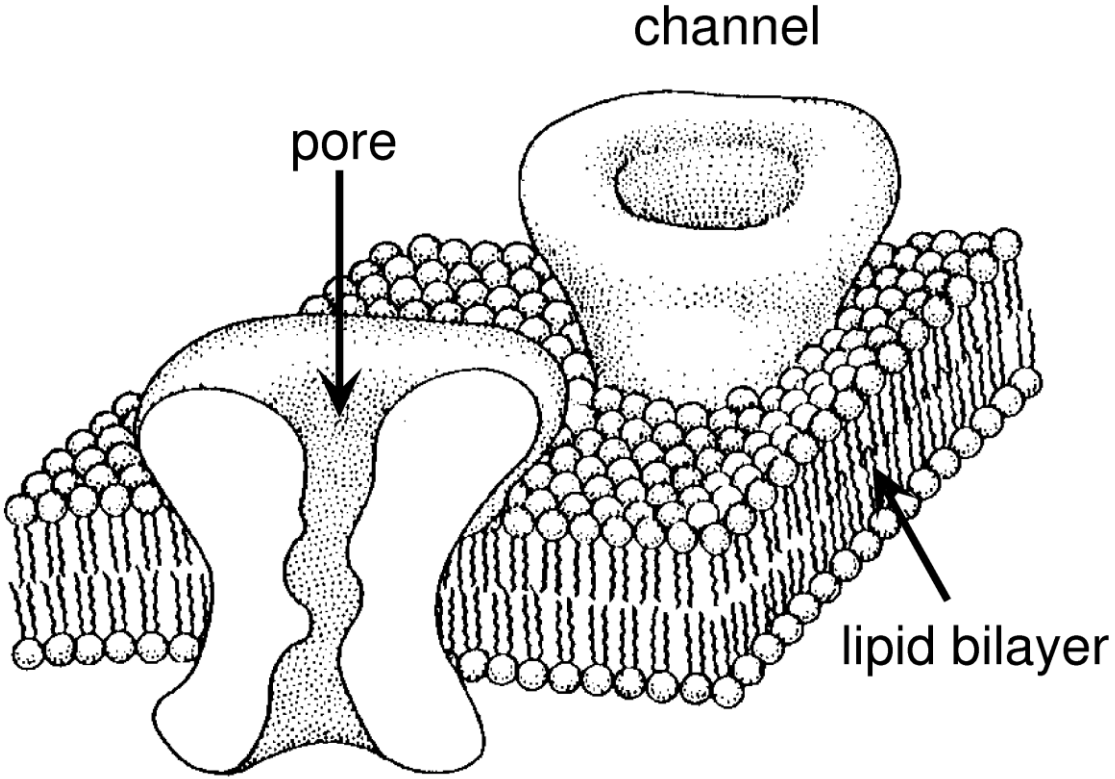
Wiring neurons together

- synapses
- long term plasticity
- short term plasticity

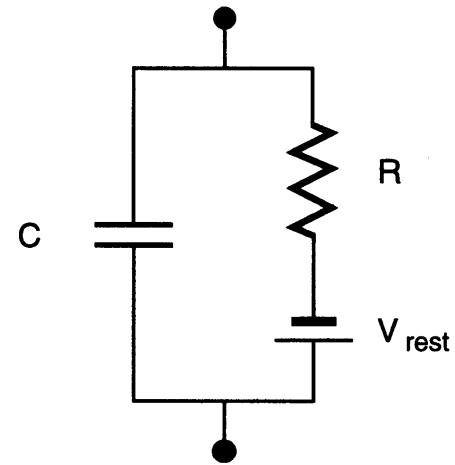
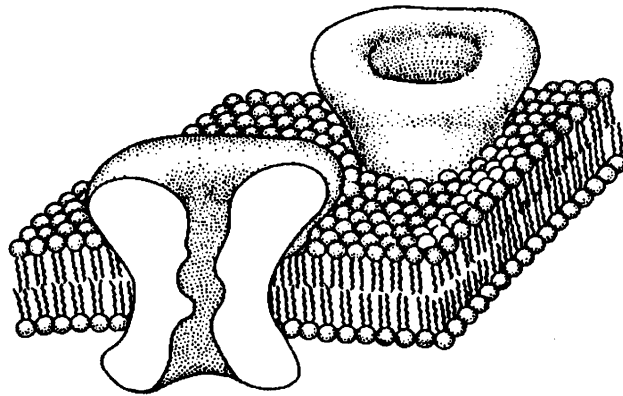
Equivalent circuit model



Membrane patch



The passive membrane



Ohm's law: $V = I_R R$

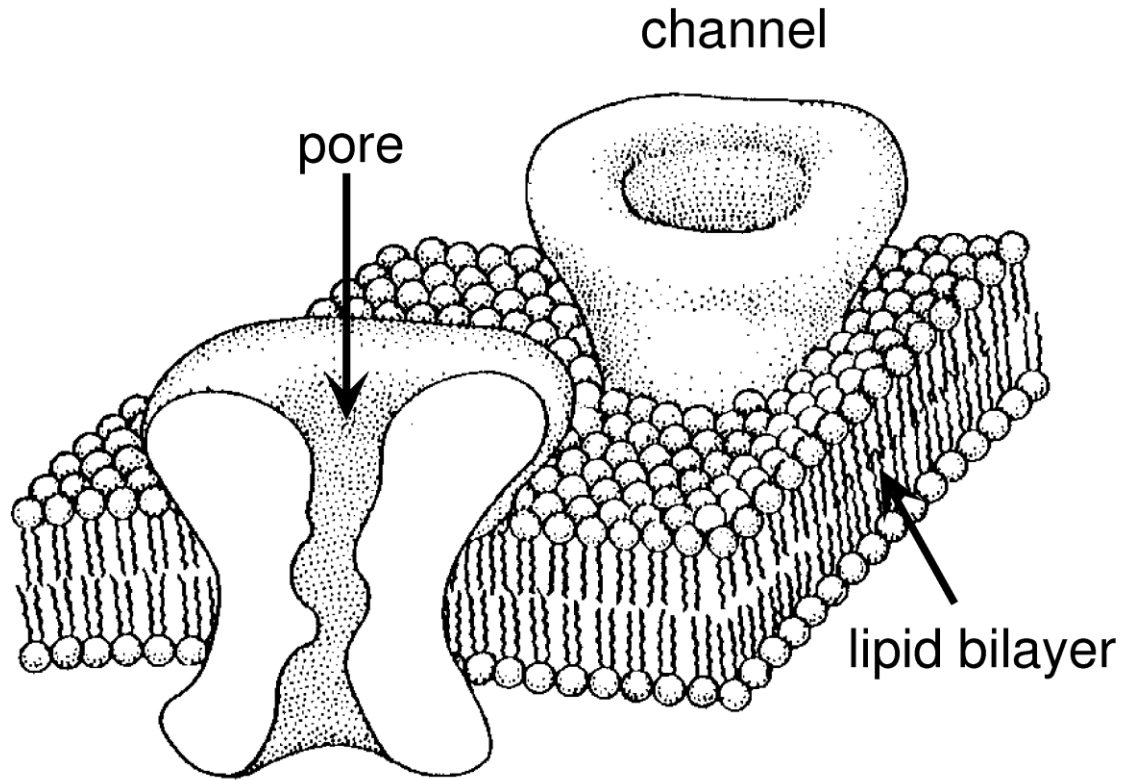
Capacitor: $C = Q/V$

$$I_C = C \frac{dV}{dt}$$

Kirchhoff: $I_R + I_C + I_{\text{ext}} = 0$

$$C \frac{dV}{dt} = -\frac{V}{R} - I_{\text{ext}}$$

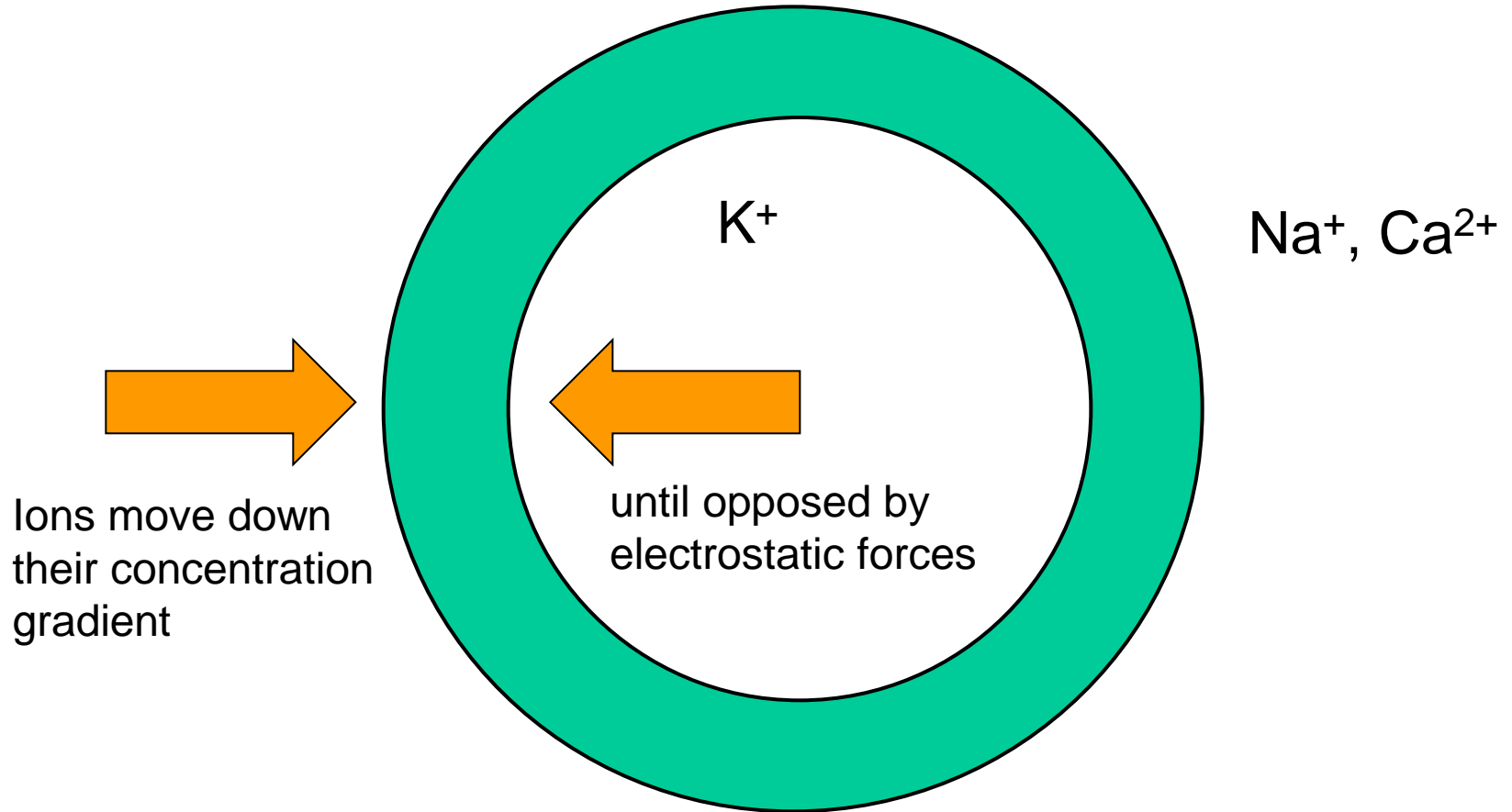
Movement of ions through ion channels



Energetics: $qV \sim k_B T$

$V \sim 25\text{mV}$

The equilibrium potential



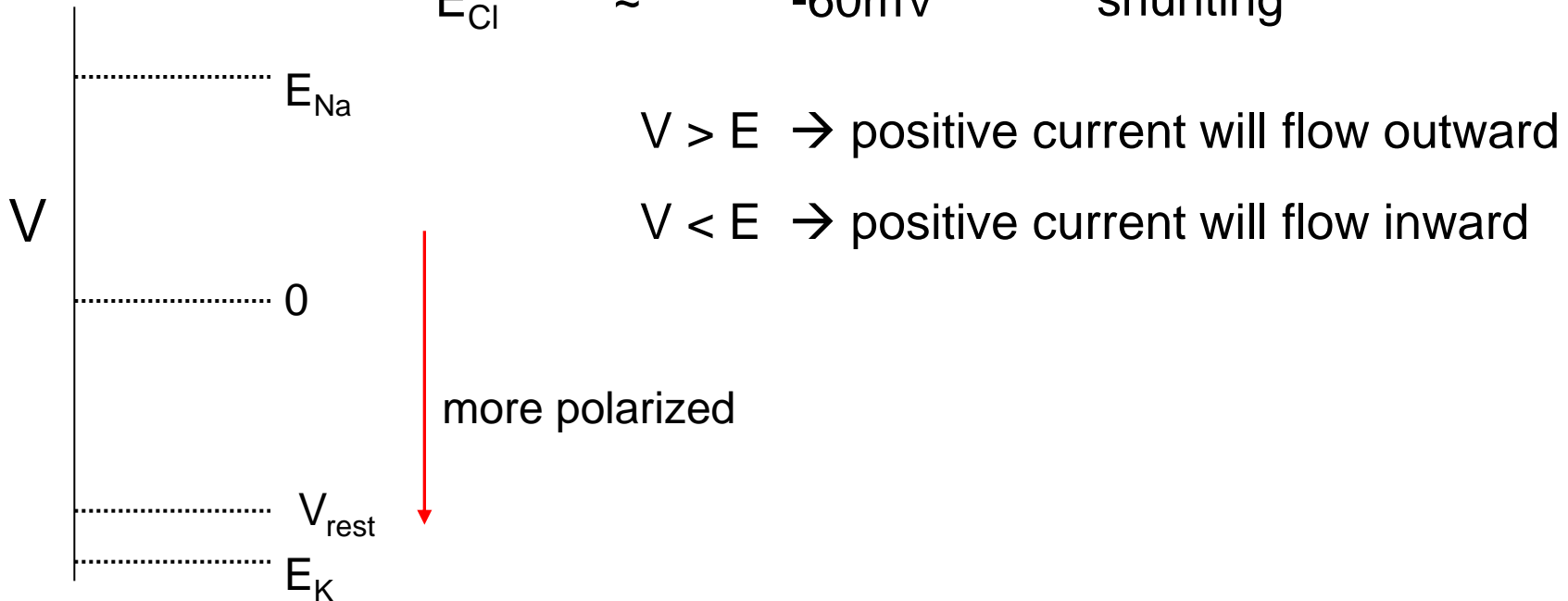
Nernst:
$$E = \frac{k_B T}{zq} \ln \frac{[\text{inside}]}{[\text{outside}]}$$

Each ion type travels through independently

Different ion channels have associated *conductances*.

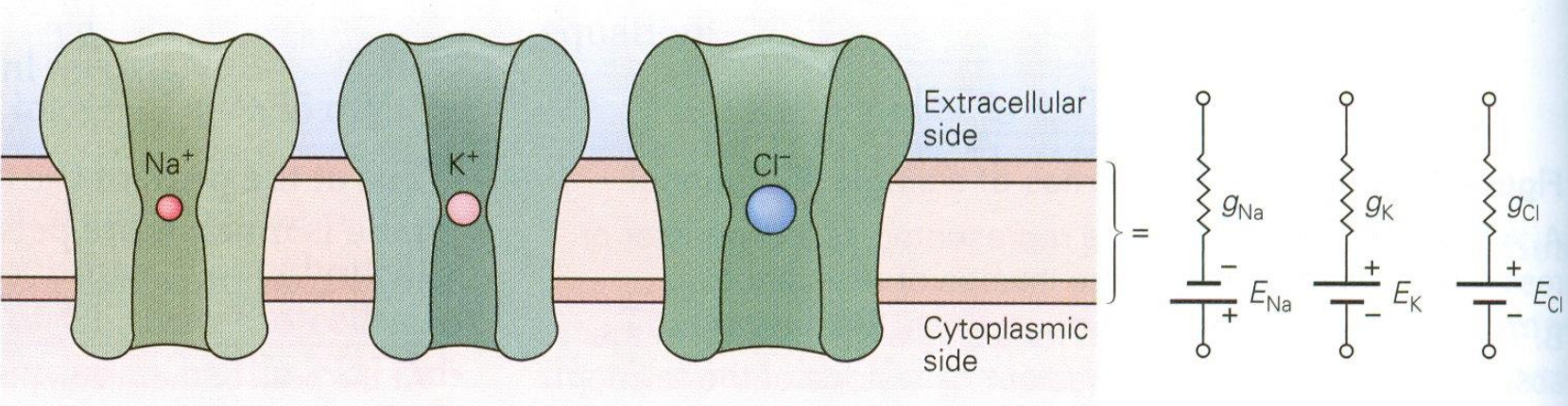
A given conductance tends to move the membrane potential toward the equilibrium potential for that ion

E_{Na}	~	50mV	depolarizing
E_{Ca}	~	150mV	depolarizing
E_K	~	-80mV	hyperpolarizing
E_{Cl}	~	-60mV	shunting

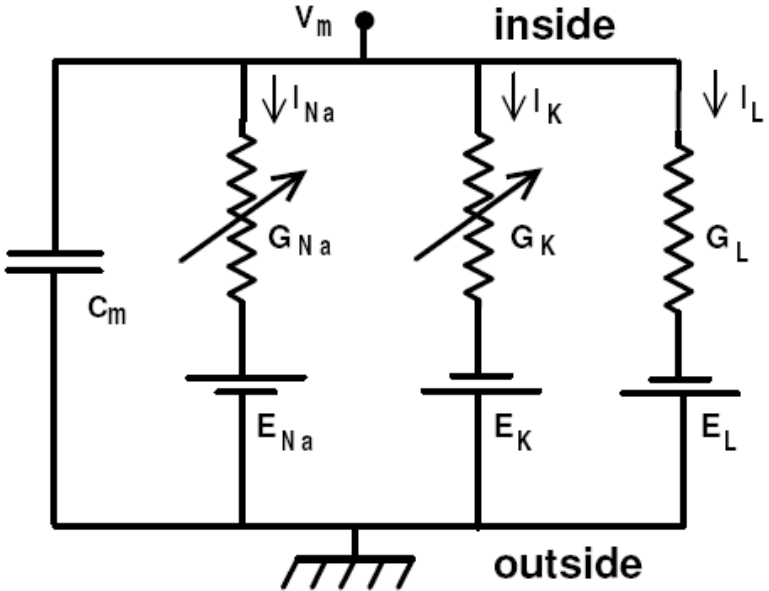


Parallel paths for ions to cross membrane

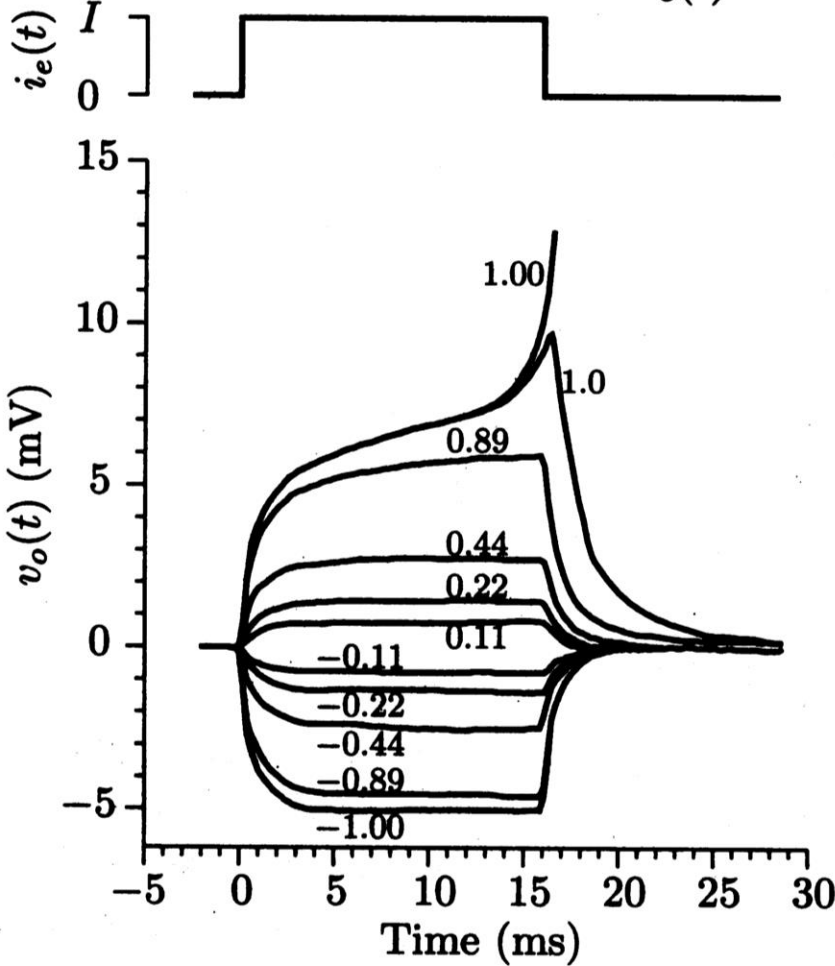
Several I - V curves in parallel:



New equivalent circuit:

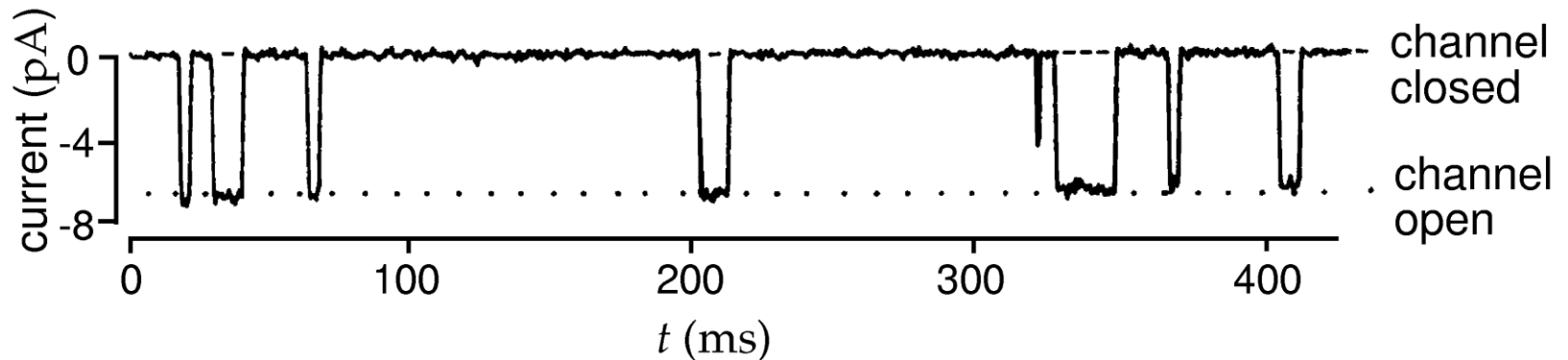
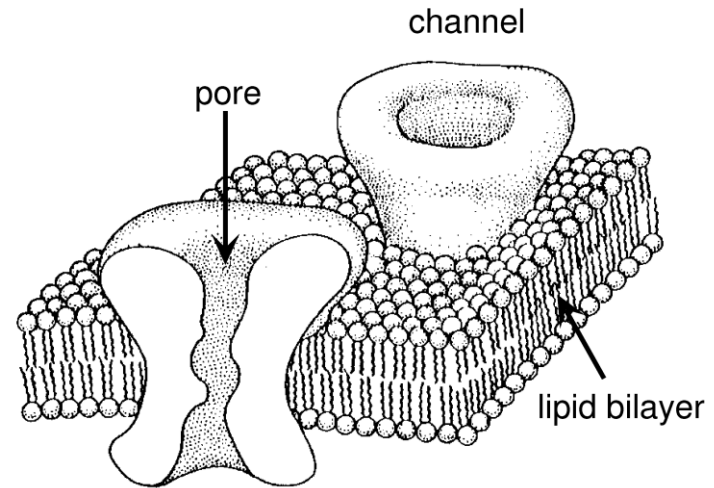


Neurons are excitable



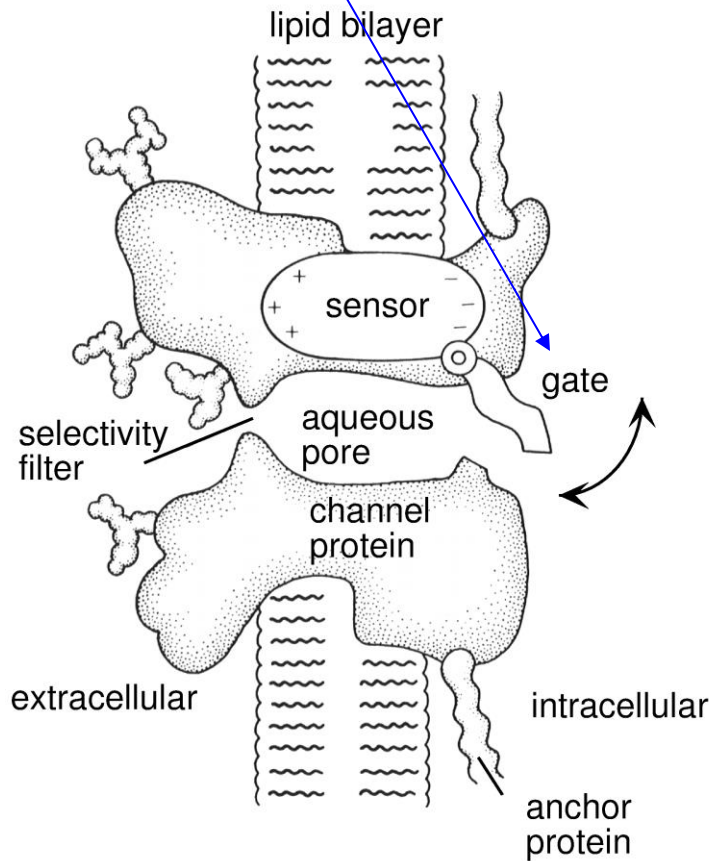
Excitability arises from ion channel nonlinearity

- Voltage dependent
- transmitter dependent (synaptic)
- Ca dependent



The ion channel is a cool molecular machine

K channel: open probability increases when depolarized



$$P_K \sim n^4$$

n describes a subunit

n is open probability
 $1 - n$ is closed probability

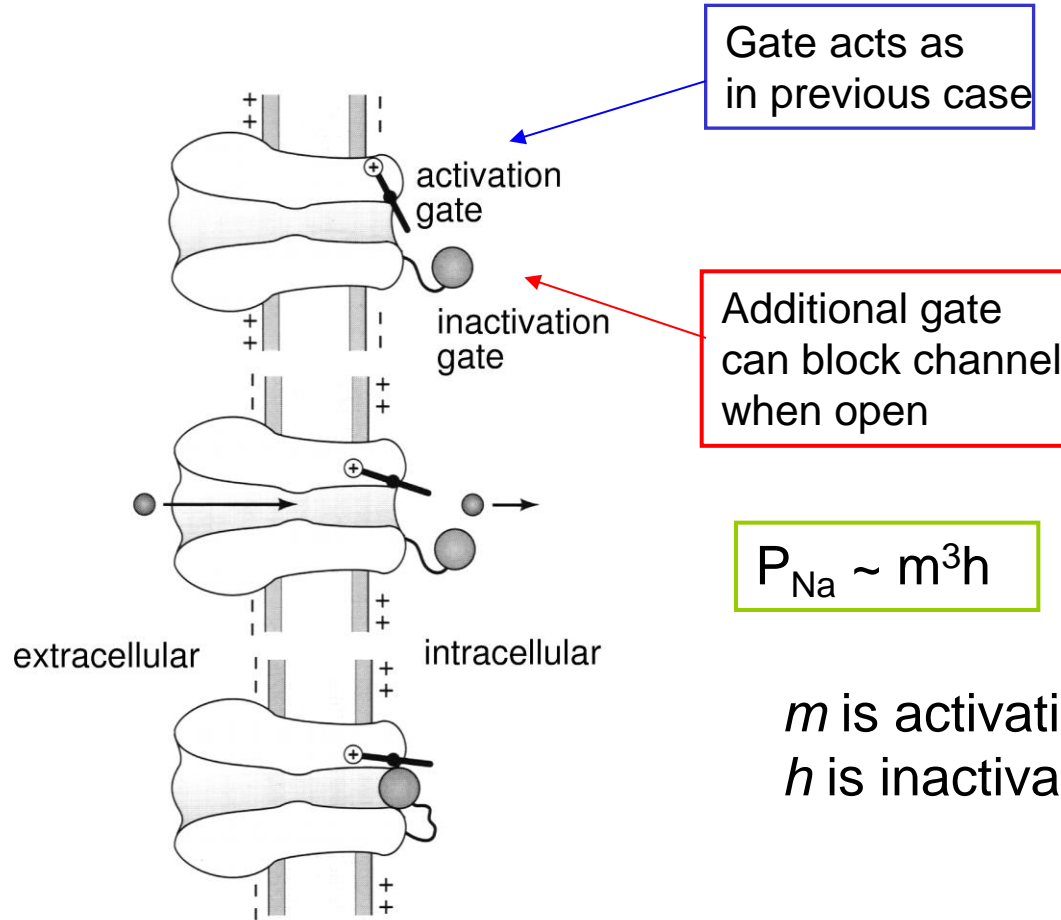
Transitions between states occur at voltage dependent rates



$$\frac{dn}{dt} = \alpha_n(V)(1 - n) - \beta_n(V)n$$

Persistent conductance

Transient conductances



m is activation variable
 h is inactivation variable

m and h have opposite voltage dependences:
depolarization increases m , activation
hyperpolarization increases h , deinactivation

Dynamics of activation and inactivation

$$\frac{dn}{dt} = \alpha_n(V)(1 - n) - \beta_n(V)n$$

$$\frac{dm}{dt} = \alpha_m(V)(1 - m) - \beta_m(V)m$$

$$\frac{dh}{dt} = \alpha_h(V)(1 - h) - \beta_h(V)h$$

We can rewrite:

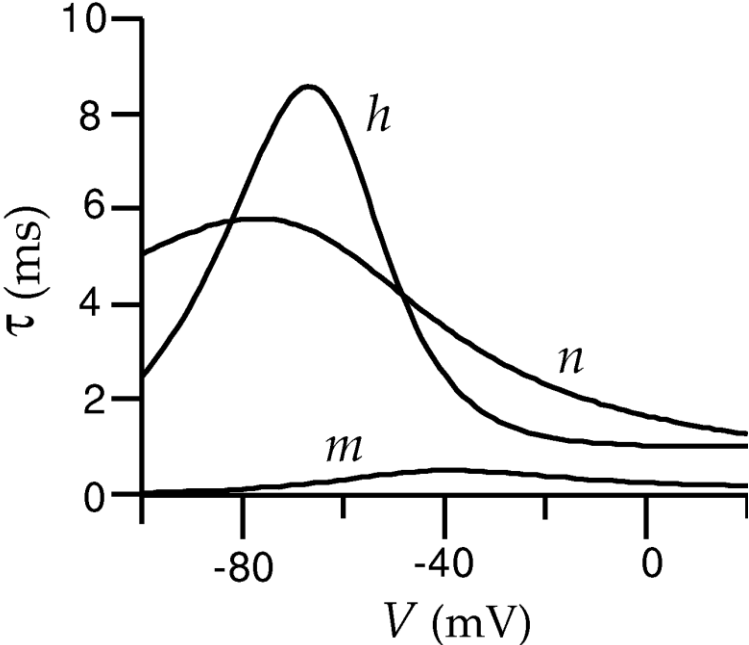
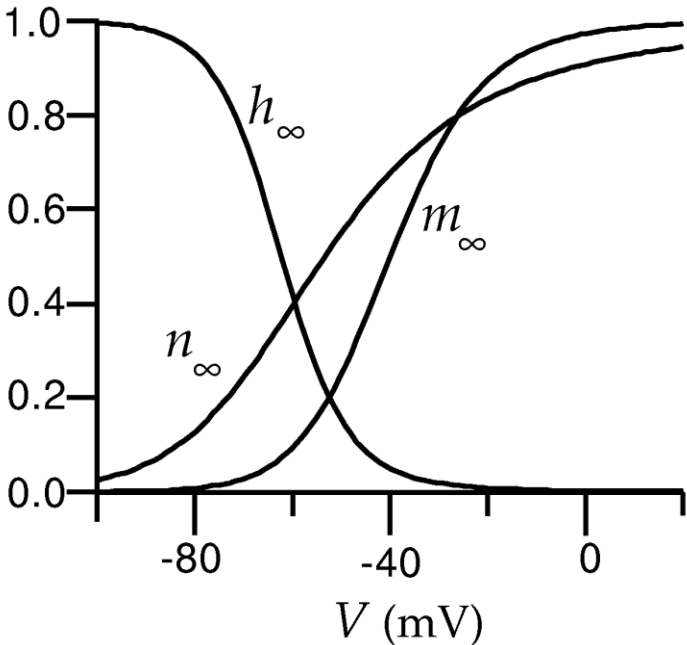
$$\tau_n(V) \frac{dn}{dt} = n_\infty(V) - n$$

where

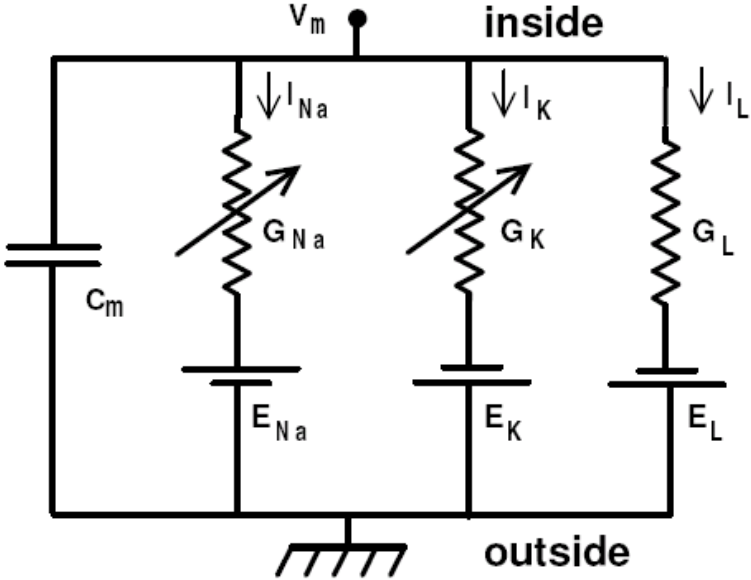
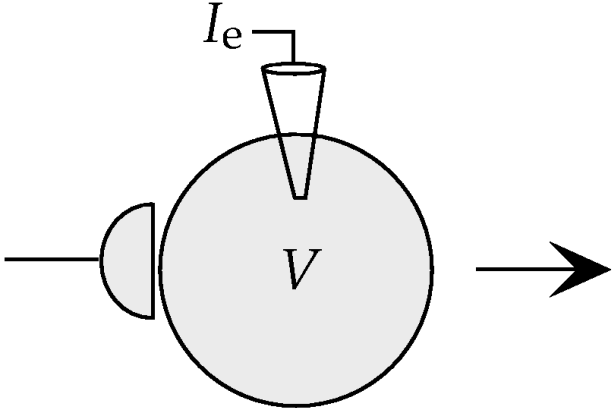
$$\tau_n(V) = \frac{1}{\alpha_n(V) + \beta_n(V)}$$

$$n_\infty(V) = \frac{\alpha_n(V)}{\alpha_n(V) + \beta_n(V)}$$

Dynamics of activation and inactivation



Putting it together



Ohm's law: $V = IR$ and Kirchoff's law

$$-C_m \frac{dV}{dt} = \sum_i g_i (V - E_i) + I_e$$

Capacitative current

Ionic currents

Externally applied current

The Hodgkin-Huxley equation

$$C_m \frac{dV}{dt} = - \sum_i g_i (V - E_i) - I_e$$

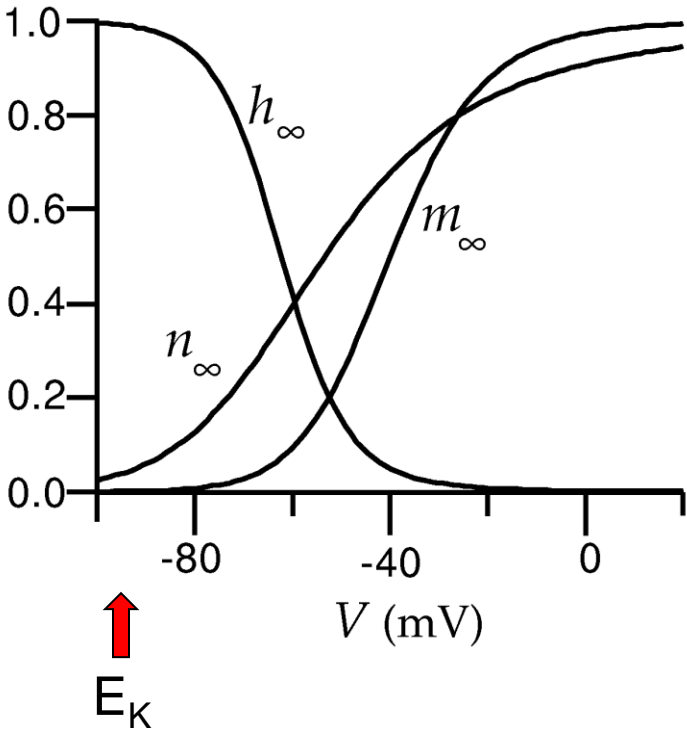
$$-C_m \frac{dV}{dt} = g_L (V - E_L) + \bar{g}_K n^4 (V - E_K) + \bar{g}_{Na} m^3 h (V - E_{Na})$$

$$\frac{dn}{dt} = \alpha_n(V)(1 - n) - \beta_n(V)n$$

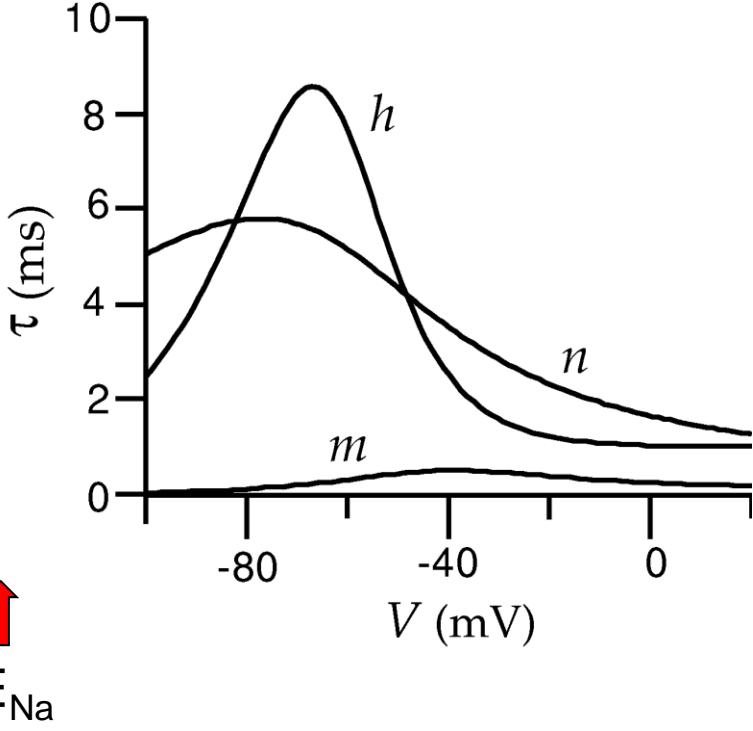
$$\frac{dm}{dt} = \alpha_m(V)(1 - m) - \beta_m(V)m$$

$$\frac{dh}{dt} = \alpha_h(V)(1 - h) - \beta_h(V)h$$

Anatomy of a spike

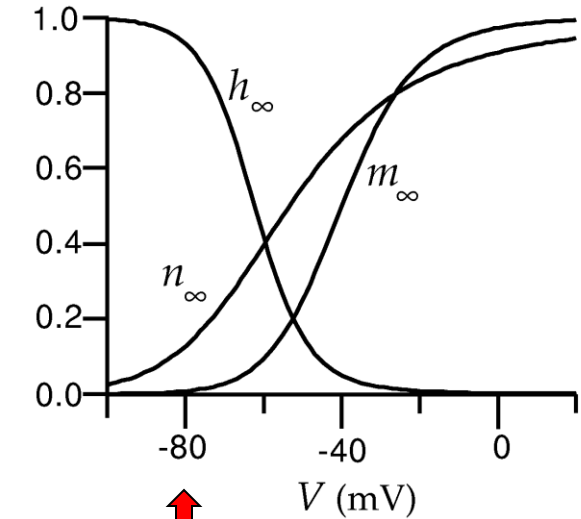


Na ~ m^3h



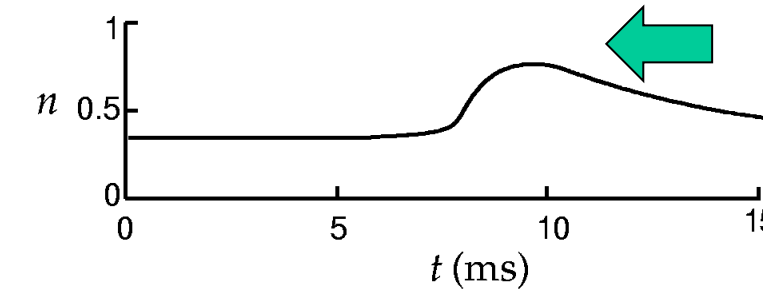
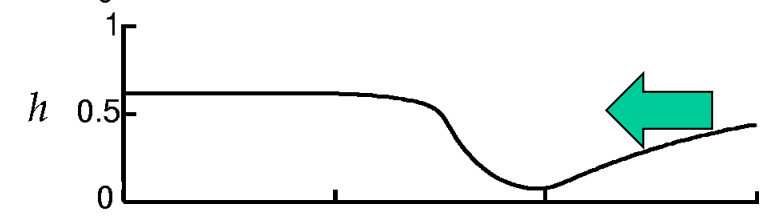
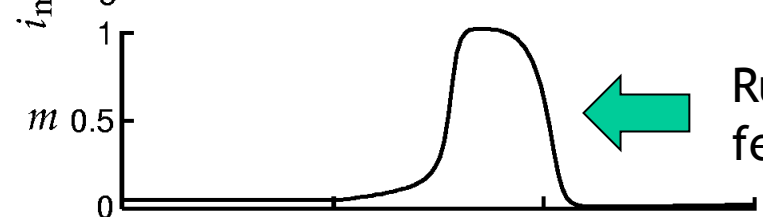
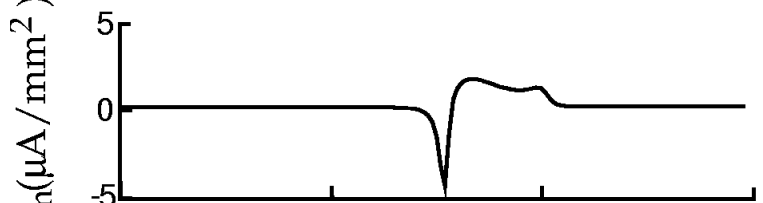
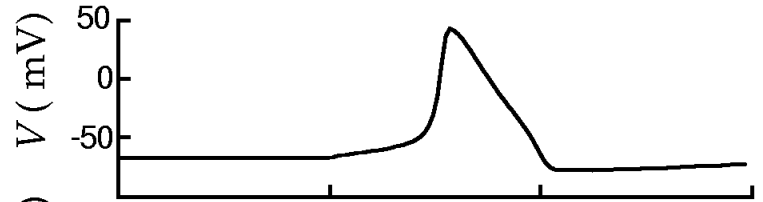
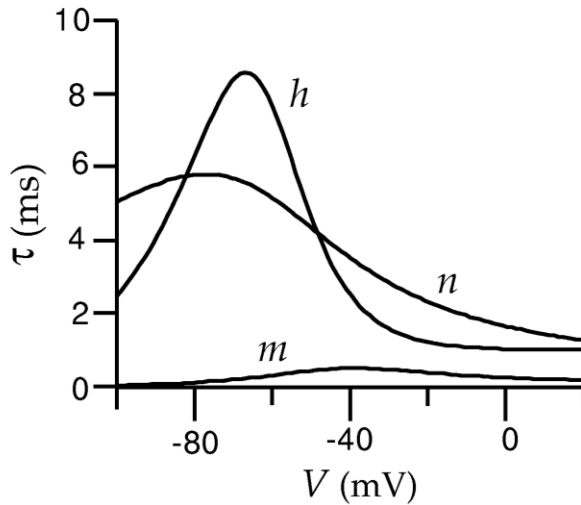
K ~ m^3h

Anatomy of a spike



↑
 E_K

↑
 E_{Na}



Where to from here?

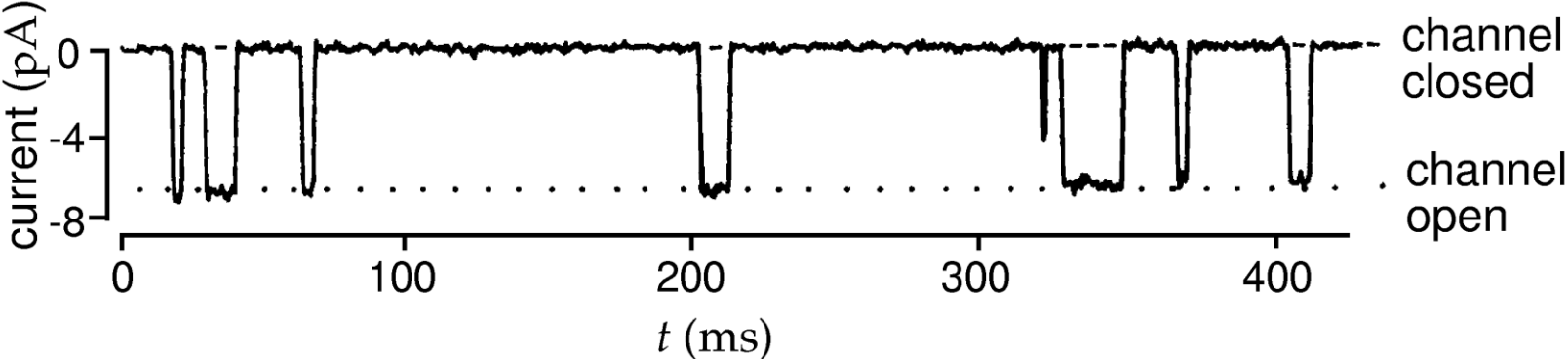
Hodgkin-Huxley

```
graph TD; A[Hodgkin-Huxley] --> B[Biophysical realism  
Molecular considerations  
Geometry]; A --> C[Simplified models  
Analytical tractability];
```

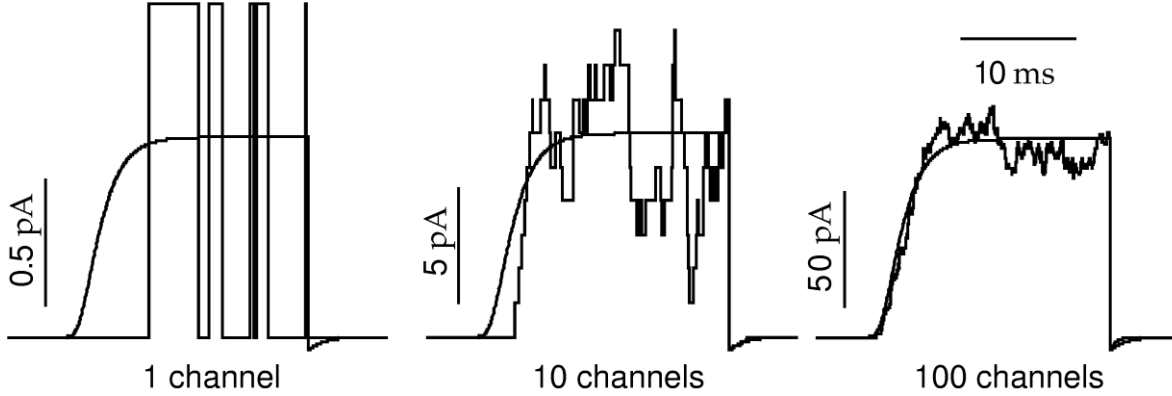
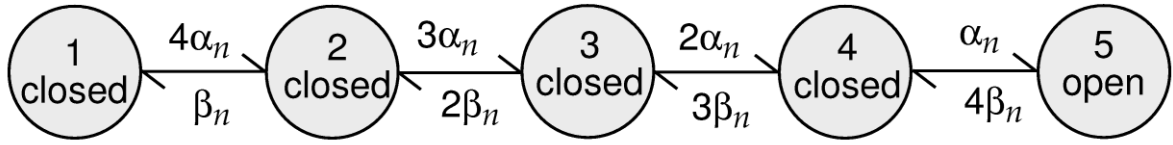
Biophysical realism
Molecular considerations
Geometry

Simplified models
Analytical tractability

Ion channel stochasticity

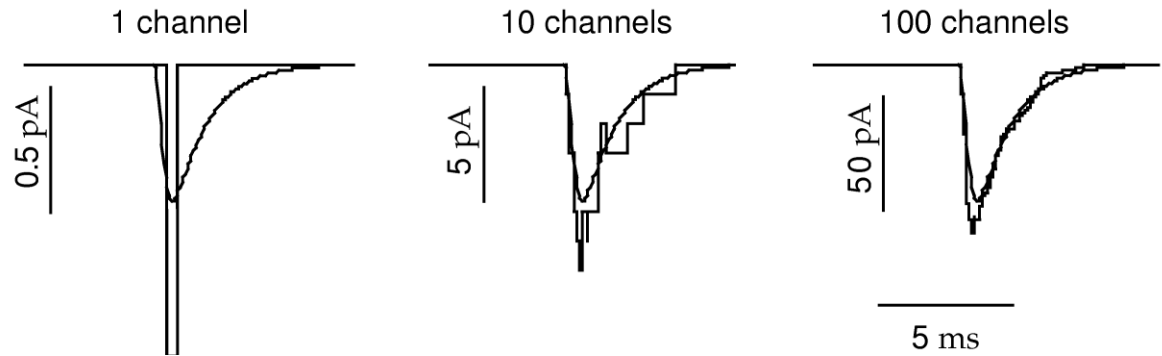
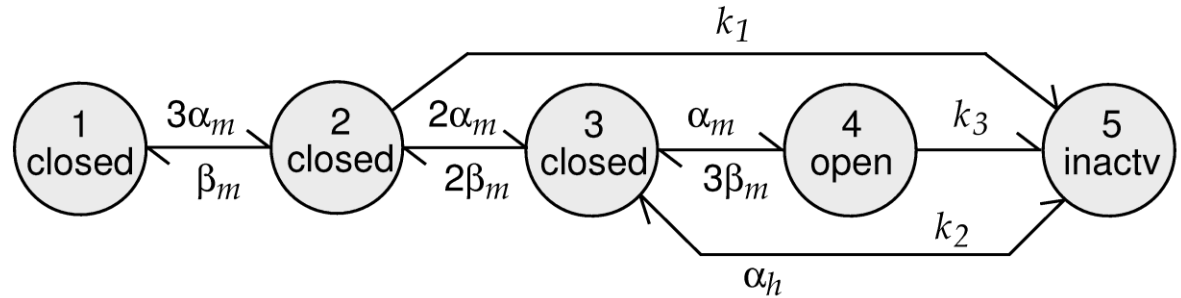
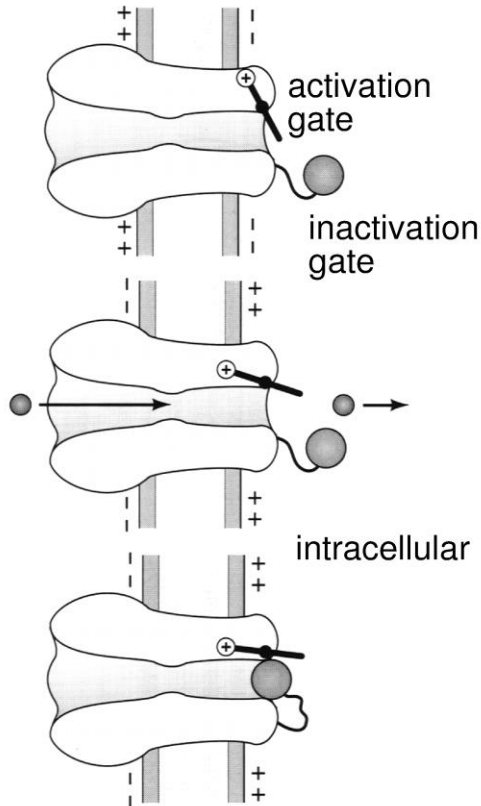


Microscopic models for ion channel fluctuations



approach to macroscopic description

Transient conductances



Different from the continuous model:

interdependence between inactivation and activation

transitions to inactivation state 5 can occur only from 2,3 and 4

k_1 , k_2 , k_3 are *constant*, not voltage dependent

The integrate-and-fire neuron

Like a passive membrane:

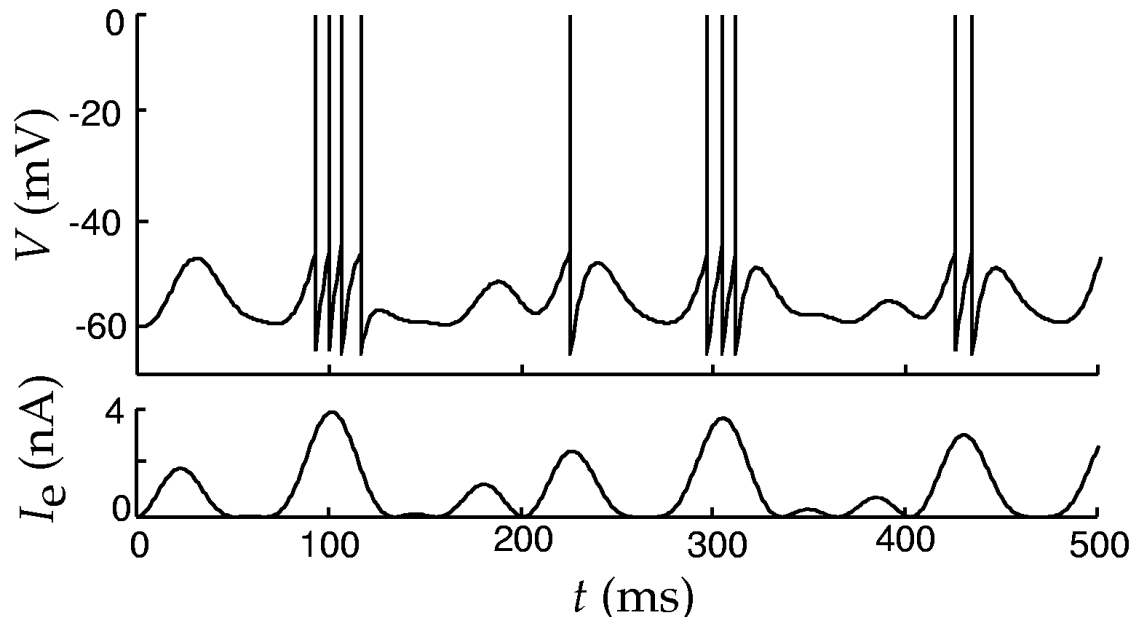
$$C_m \frac{dV}{dt} = -g_L(V - E_i) - I_e$$

but with the additional rule that

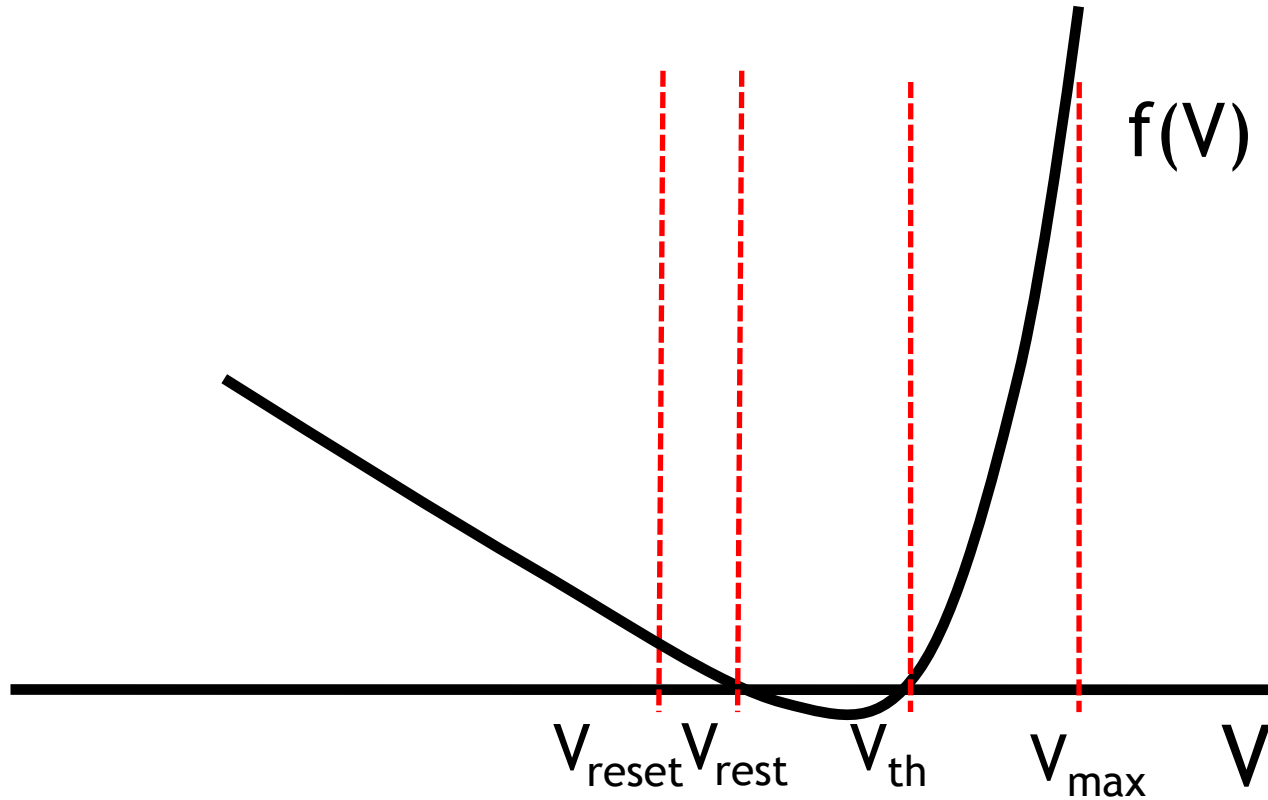
when $V \rightarrow V_T$, a spike is fired

and $V \rightarrow V_{\text{reset}}$

E_L is the resting potential of the “cell”.

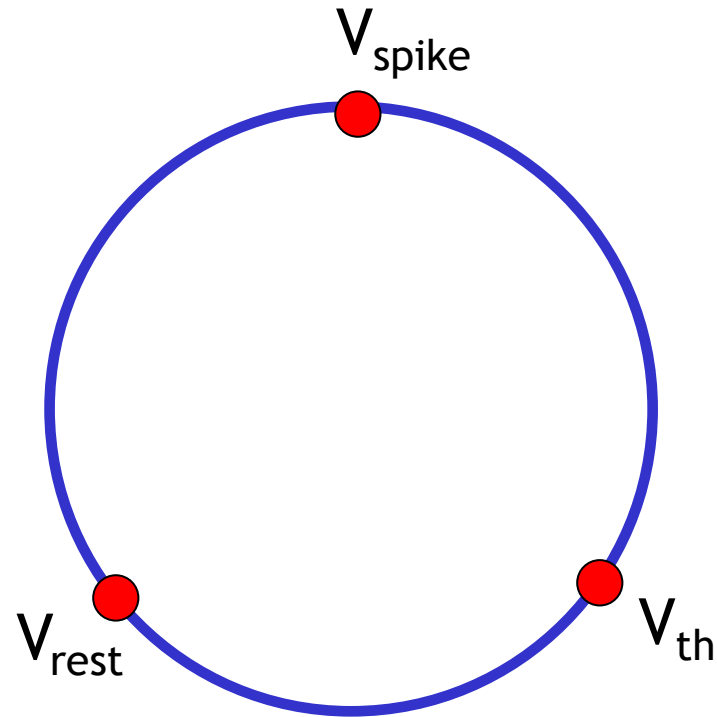


Exponential integrate-and-fire neuron



$$f(V) = -V + \exp([V - V_{th}] / \Delta)$$

The theta neuron



$$d\theta/dt = 1 - \cos \theta + (1 + \cos \theta) I(t)$$

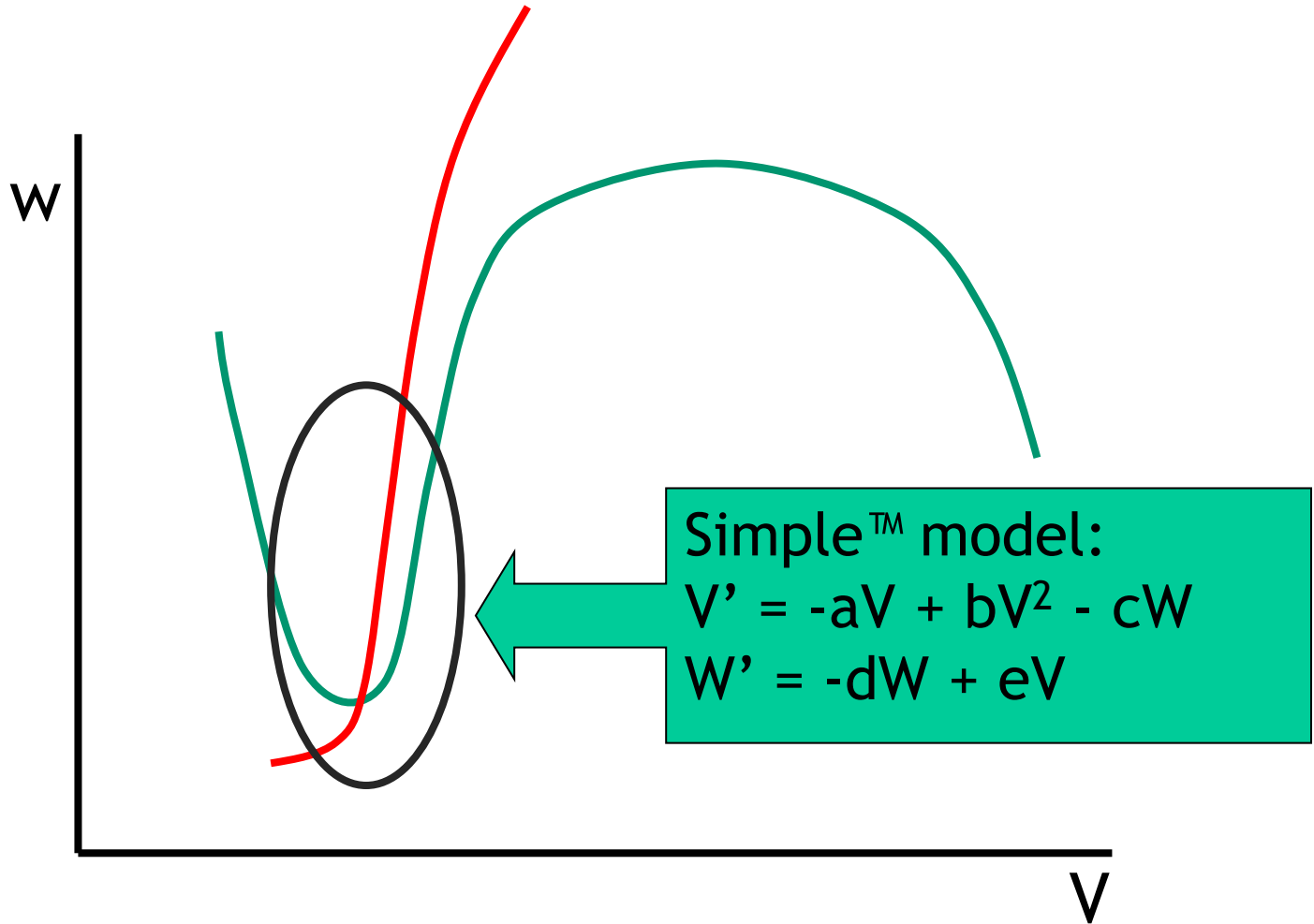
Ermentrout and Kopell

The spike response model

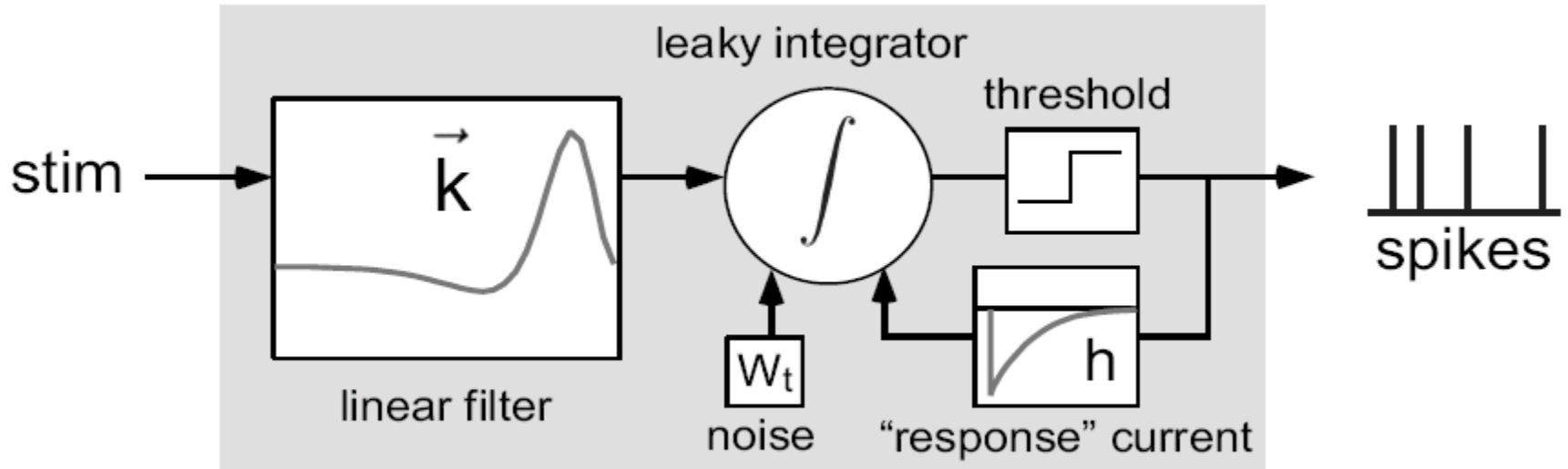
Kernel f for subthreshold response \leftarrow replaces leaky integrator
Kernel for spikes \leftarrow replaces “line”

- determine f from the linearized HH equations
- fit a threshold
- paste in the spike shape and AHP

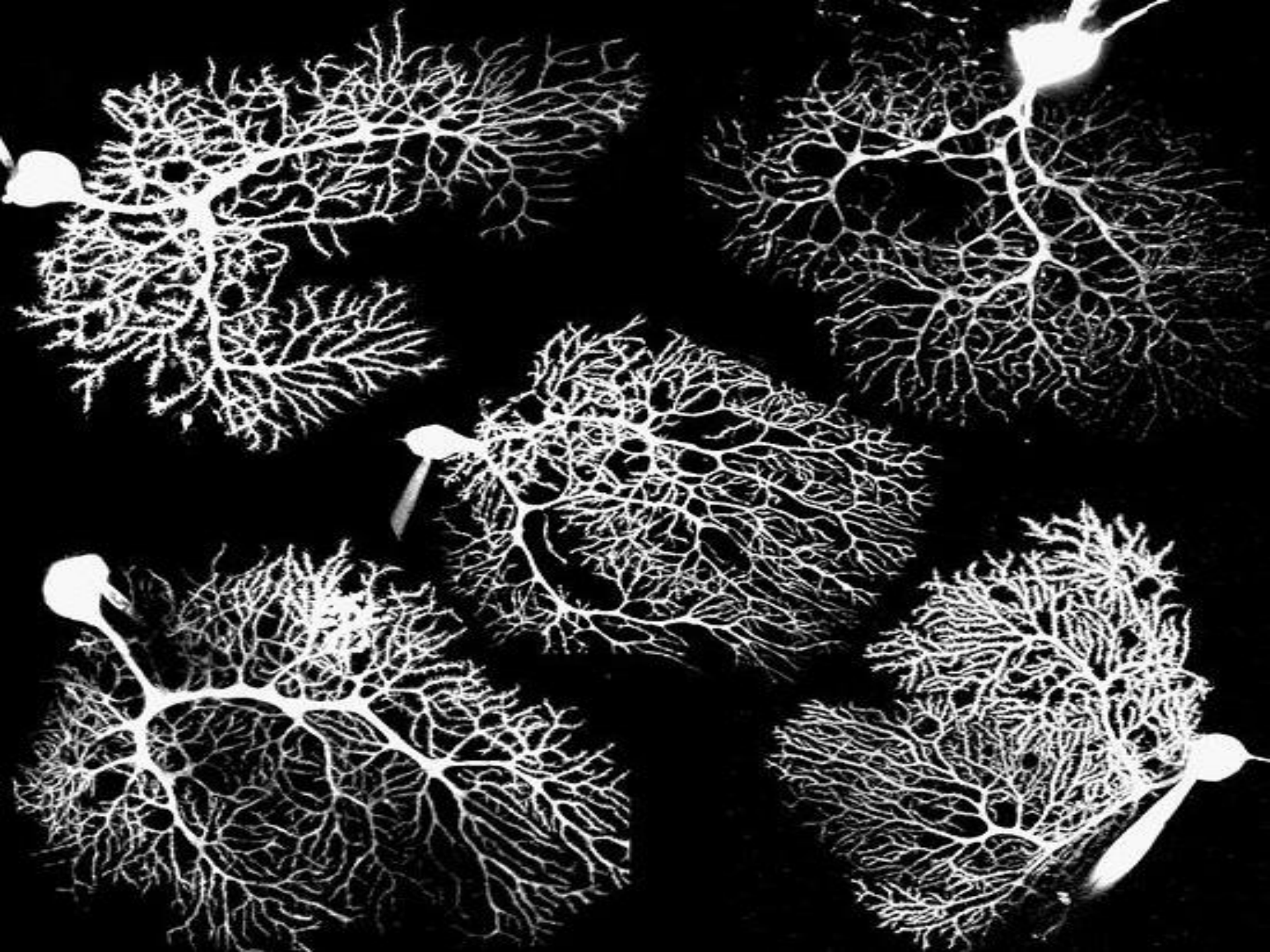
Two-dimensional models



The generalized linear model



- general definitions for k and h
- robust maximum likelihood fitting procedure



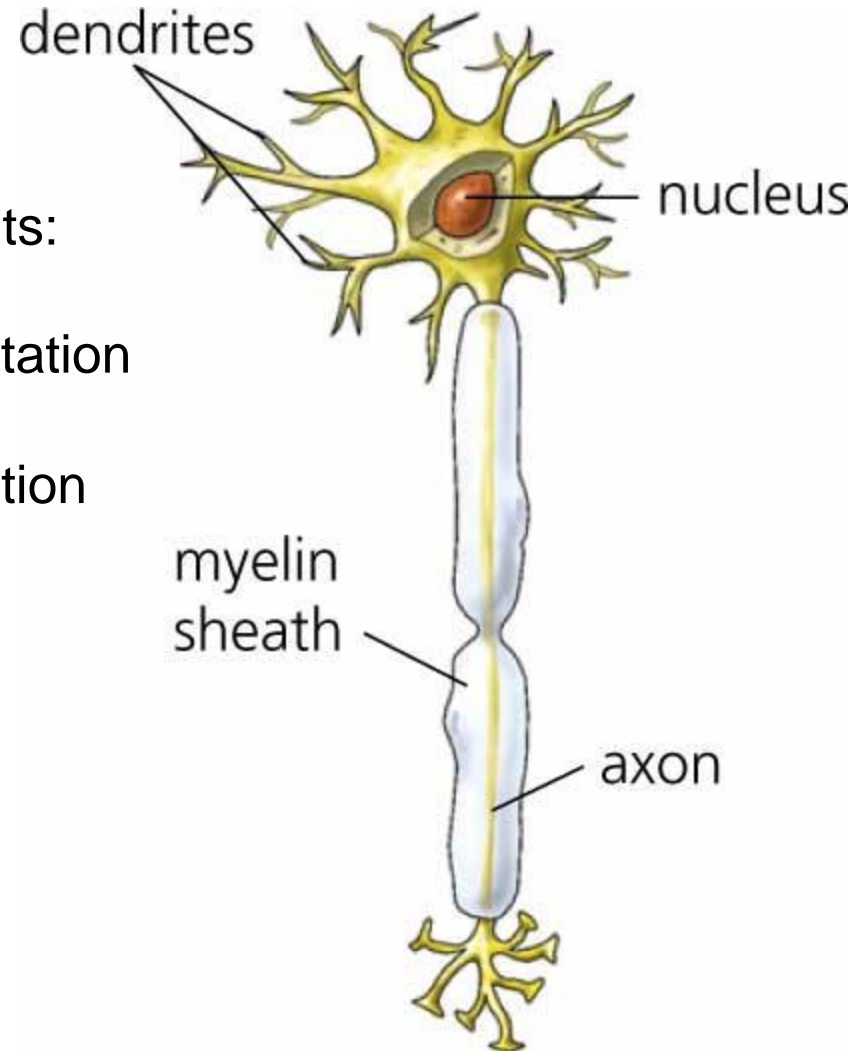
Dendritic computation

Dendrites as computational elements:

Passive contributions to computation

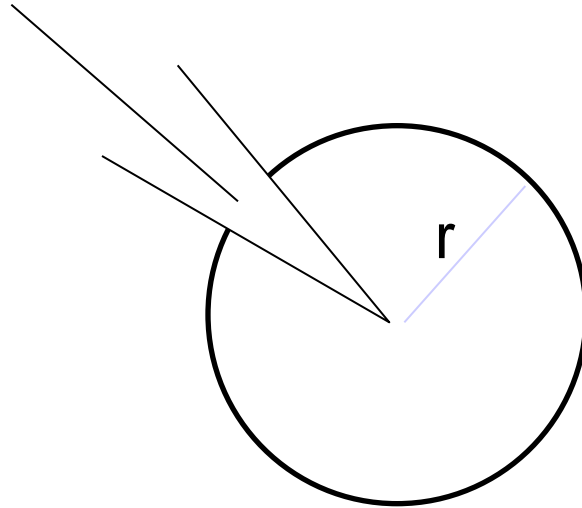
Active contributions to computation

Examples



Geometry matters

Injecting
current I_0



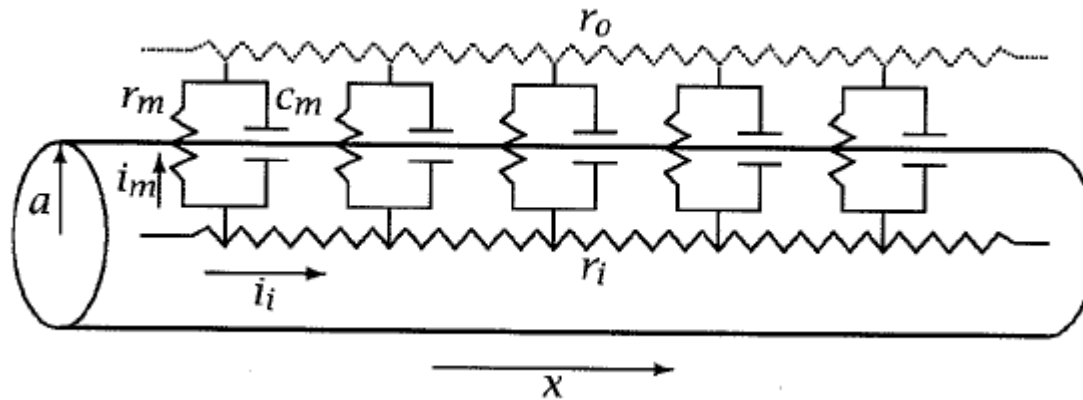
$$V_m = I_m R_m$$

Current flows uniformly out through the cell: $I_m = I_0/4\pi r^2$

Input resistance is defined as $R_N = V_m(t \rightarrow \infty)/I_0$

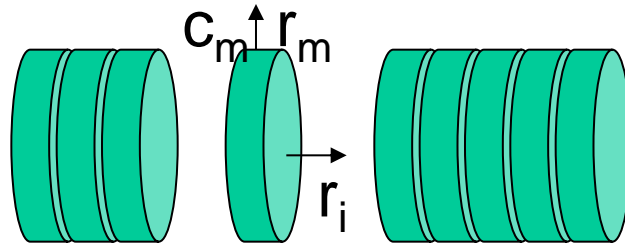
$$= R_m/4\pi r^2$$

Linear cables



r_m and r_i are the membrane and axial resistances, i.e. the resistances of a thin slice of the cylinder

Axial and membrane resistance



For a length L of membrane cable:

$$r_i \rightarrow r_i L$$

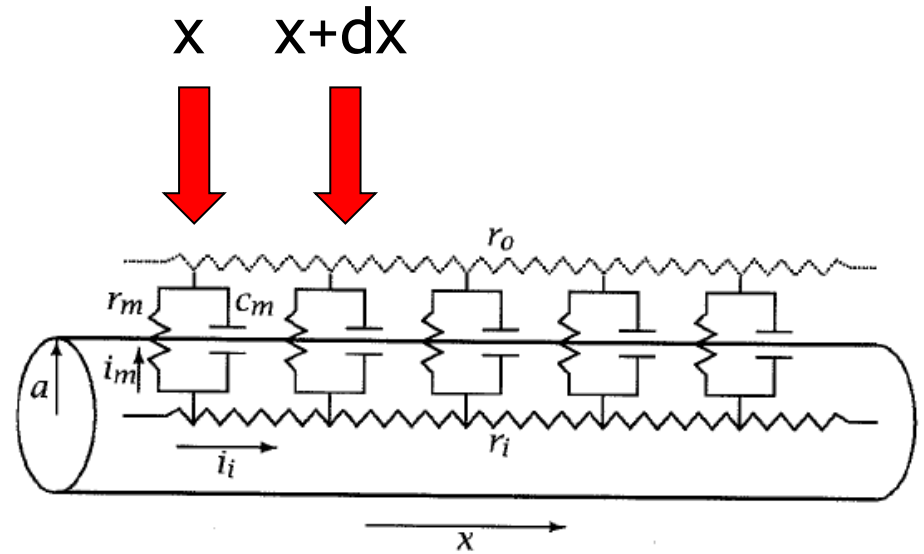
$$r_m \rightarrow r_m / L$$

$$C_m \rightarrow C_m L$$

The cable equation

$$(1) \quad \frac{\partial V_m}{\partial x} = -r_i i_i$$

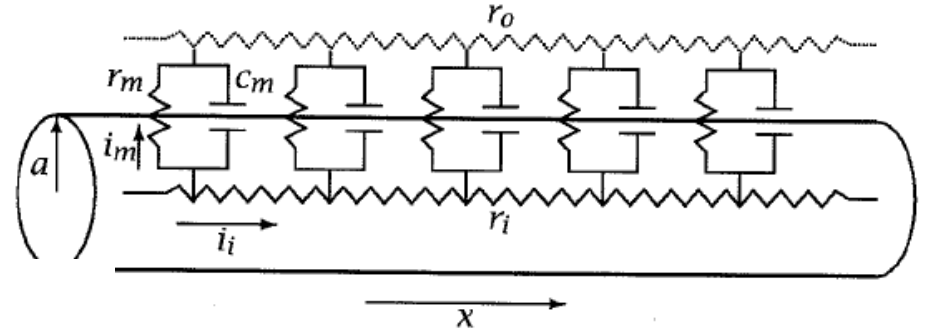
$$(2) \quad \frac{\partial i_i(x)}{\partial x} = -i_m$$



The cable equation

$$(1) \quad \frac{\partial V_m}{\partial x} = -r_i i_i$$

$$(2) \quad \frac{\partial i_i(x)}{\partial x} = -i_m$$



$$\frac{\partial}{\partial x} (1) \Rightarrow \frac{\partial^2 V_m}{\partial x^2} = -r_i \frac{\partial i_i}{\partial x} = r_i i_m.$$

$$i_m = i_C + i_{\text{ionic}} = c_m \frac{\partial V_m}{\partial t} + \frac{V_m}{r_m}$$

$$\frac{1}{r_i} \frac{\partial^2 V_m(x, t)}{\partial x^2} = c_m \frac{\partial V}{\partial t} + \frac{V_m}{r_m}.$$

or

$$\lambda^2 \frac{\partial^2 V_m}{\partial x^2} = \tau_m \frac{\partial V_m}{\partial t} + V_m$$

where

$$\tau_m = r_m c_m$$

$$\lambda = \sqrt{\frac{r_m}{r_i}}$$

Time constant

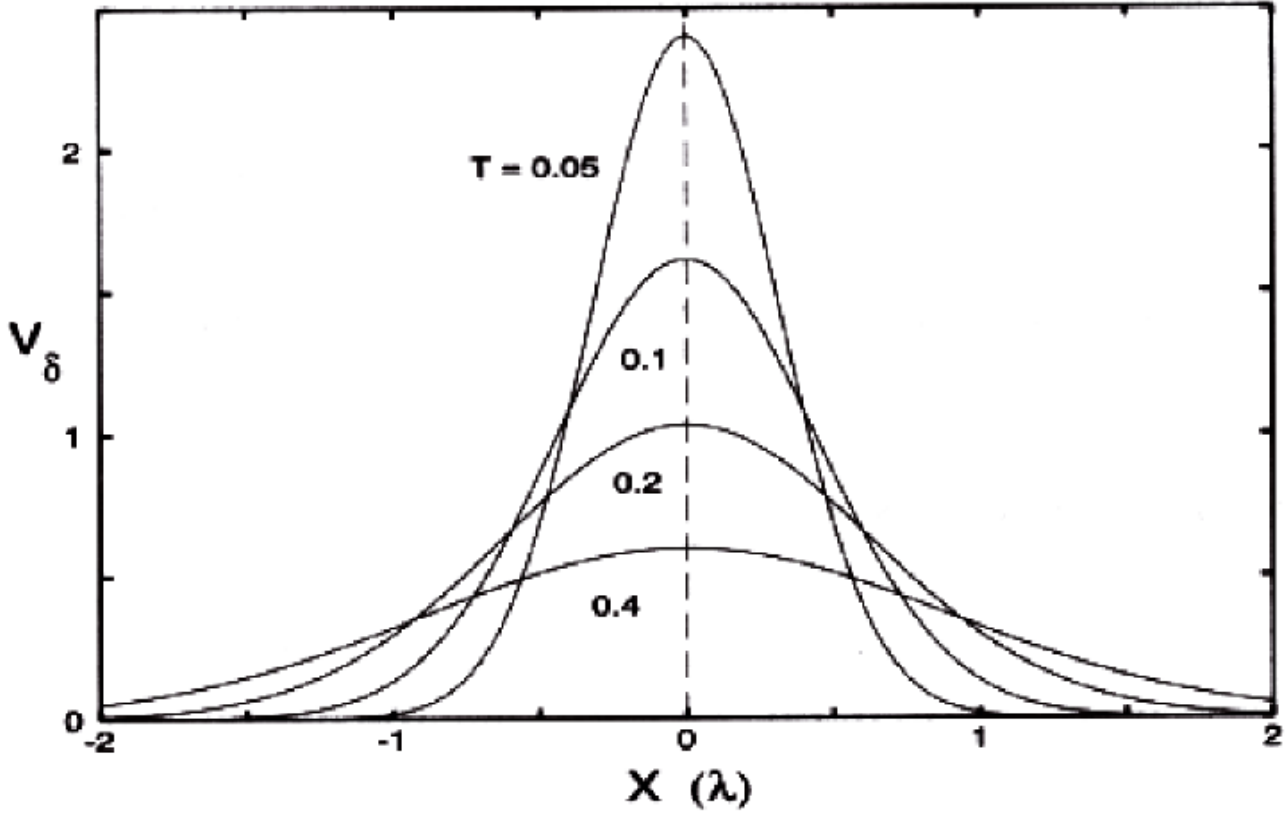
Space constant

General solution: filter and impulse response

$$V(x, t) \propto \sqrt{\frac{\tau}{4\pi\lambda^2 t}} e^{-\frac{t}{\tau} - \frac{\tau x^2}{4\lambda^2 t}}$$

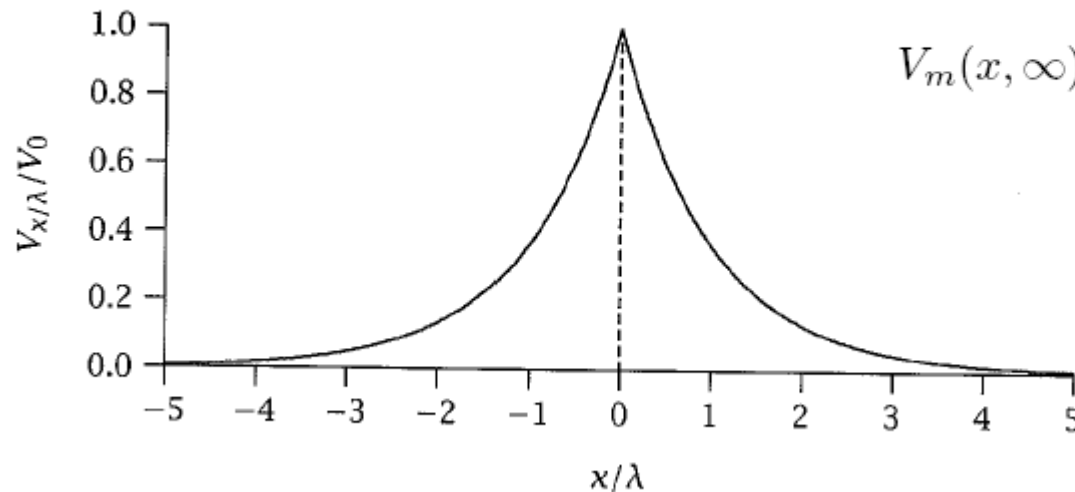
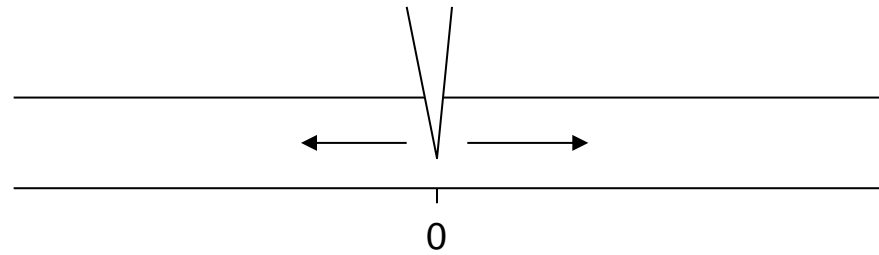
Exponential decay

Diffusive spread



Voltage decays exponentially away from source

Current injection at $x=0$, $T \rightarrow \infty$



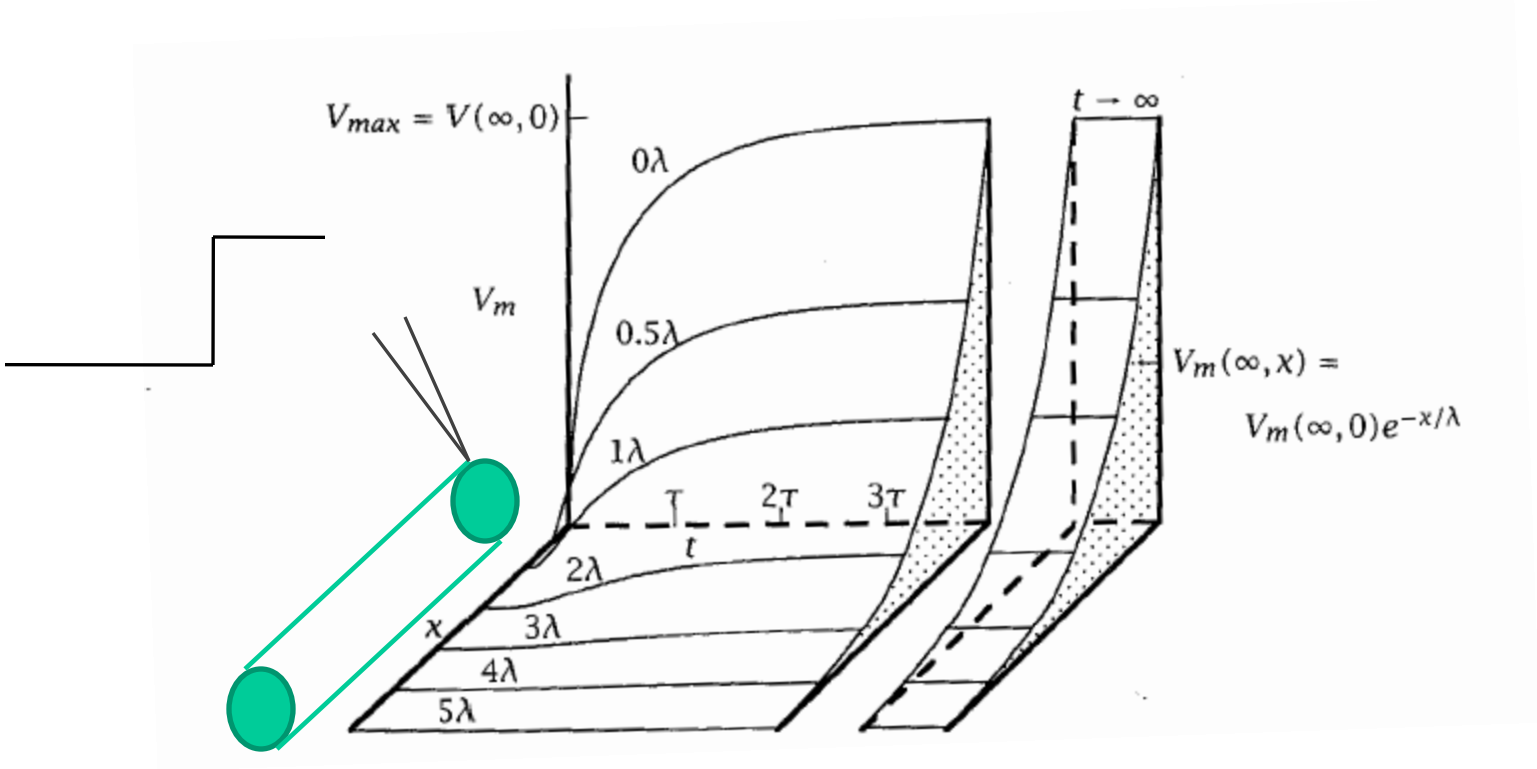
$$V_m(x, \infty) = \frac{r_i I_0 \lambda}{2} e^{-x/\lambda}$$

Properties of passive cables

→ Electrotonic length

$$\lambda = \sqrt{\frac{r_m}{r_i}}$$

Electrotonic length

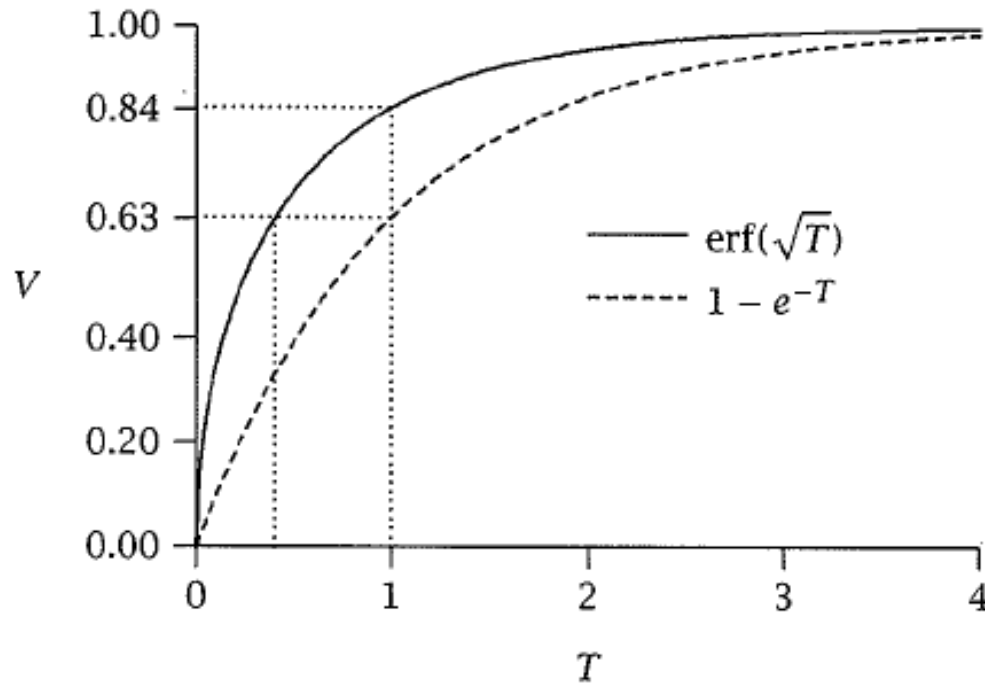


Properties of passive cables

- Electrotonic length $\lambda = \sqrt{\frac{r_m}{r_i}}$
- Current can escape through additional pathways:
speeds up decay

Voltage rise time

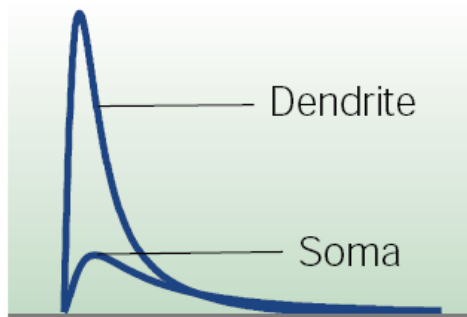
→ Current can escape through additional pathways:
speeds up decay



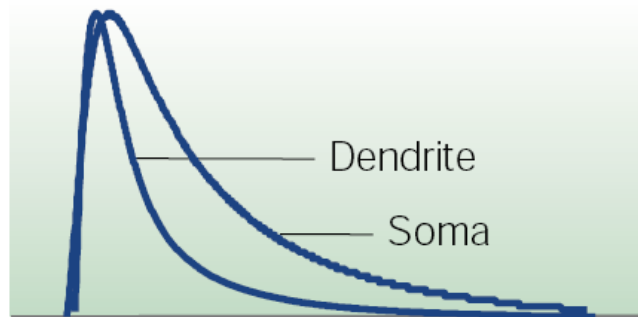
Properties of passive cables

- Electrotonic length $\lambda = \sqrt{\frac{r_m}{r_i}}$
- Current can escape through additional pathways: speeds up decay
- Cable diameter affects input resistance $R_N = \frac{\sqrt{R_m R_i / 2}}{2\pi a^{3/2}}$

: Amplitude



Time course



Properties of passive cables

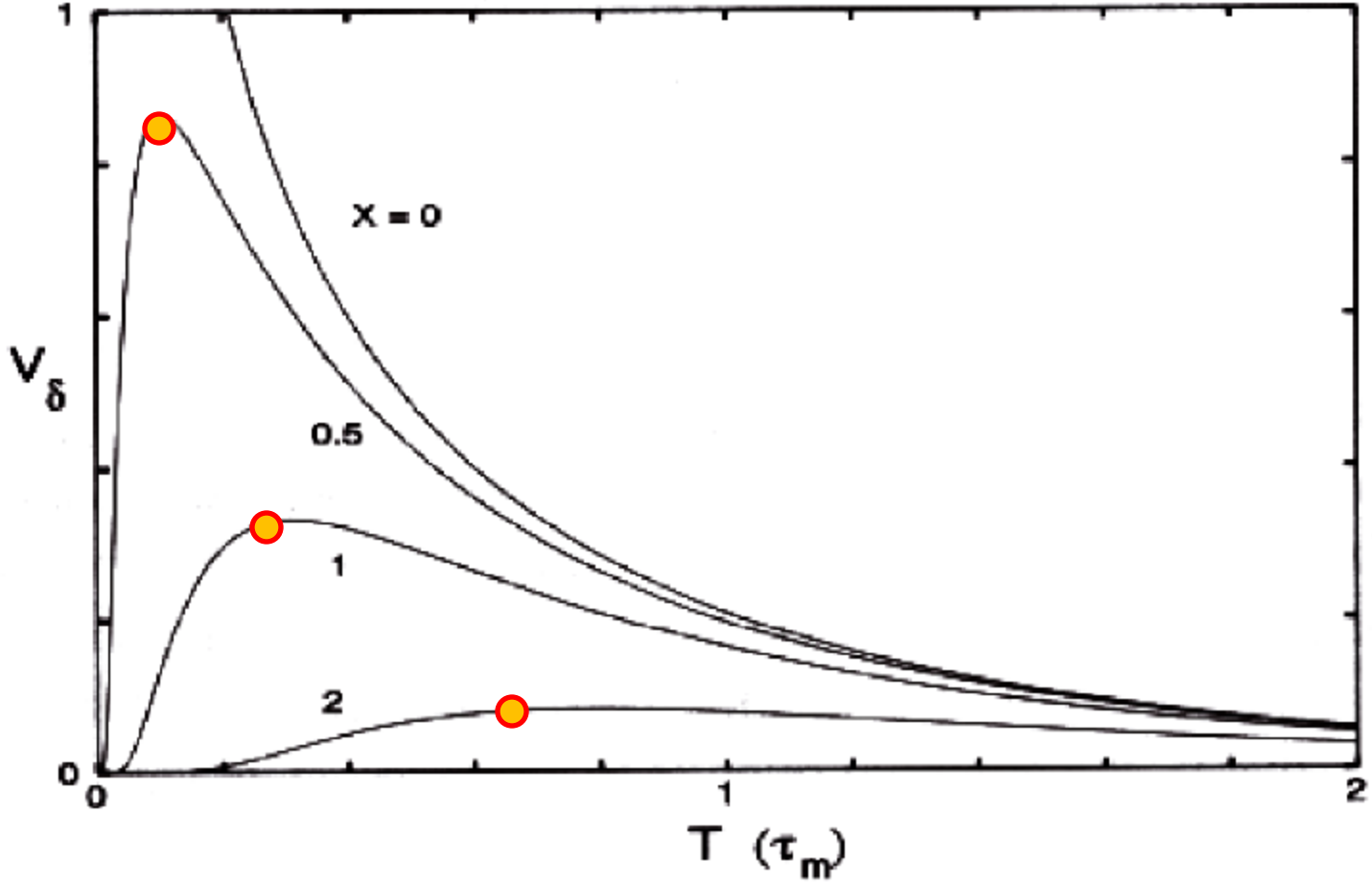
→ Electrotonic length $\lambda = \sqrt{\frac{r_m}{r_i}}$

→ Current can escape through additional pathways:
speeds up decay

→ Cable diameter affects input resistance $R_N = \frac{\sqrt{R_m R_i / 2}}{2\pi a^{3/2}}$

→ Cable diameter affects transmission velocity

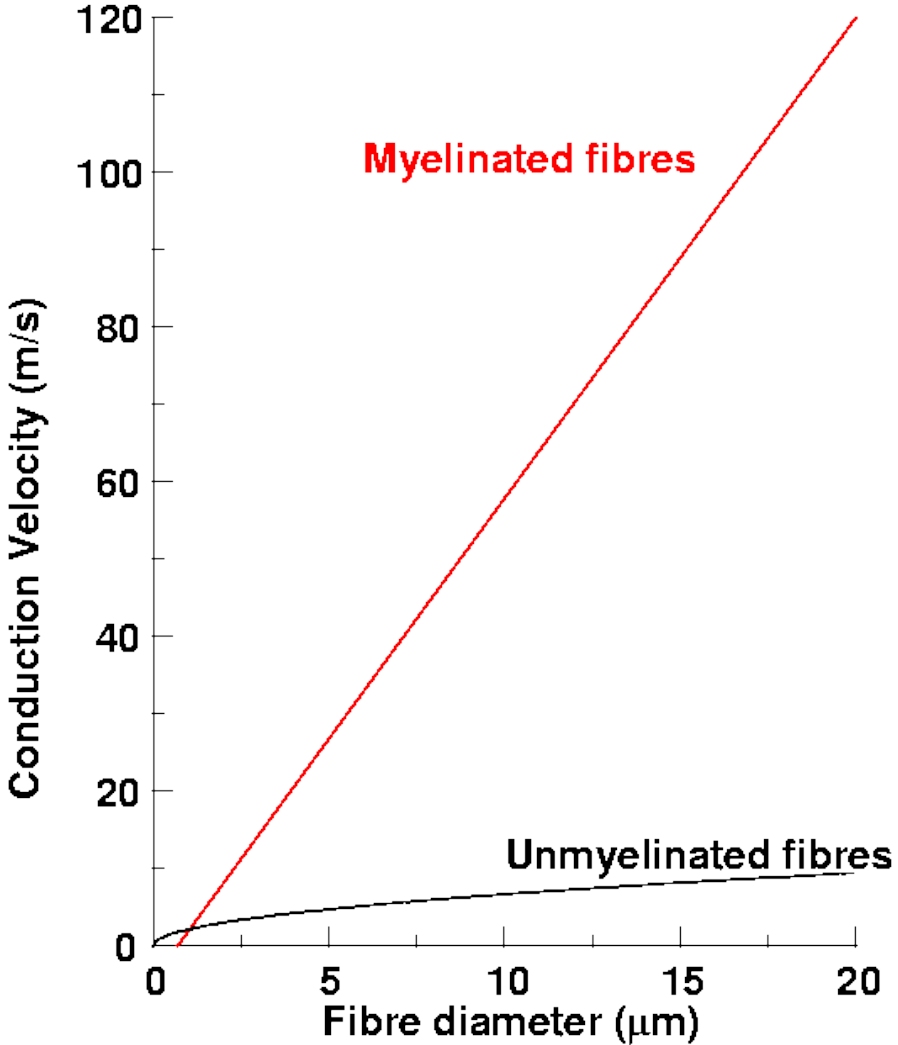
Step response: pulse travels



Conduction velocity

$$\theta = \frac{2\lambda}{\tau_m} = \sqrt{\frac{2a}{R_m R_i C_m^2}}$$

Conduction velocity

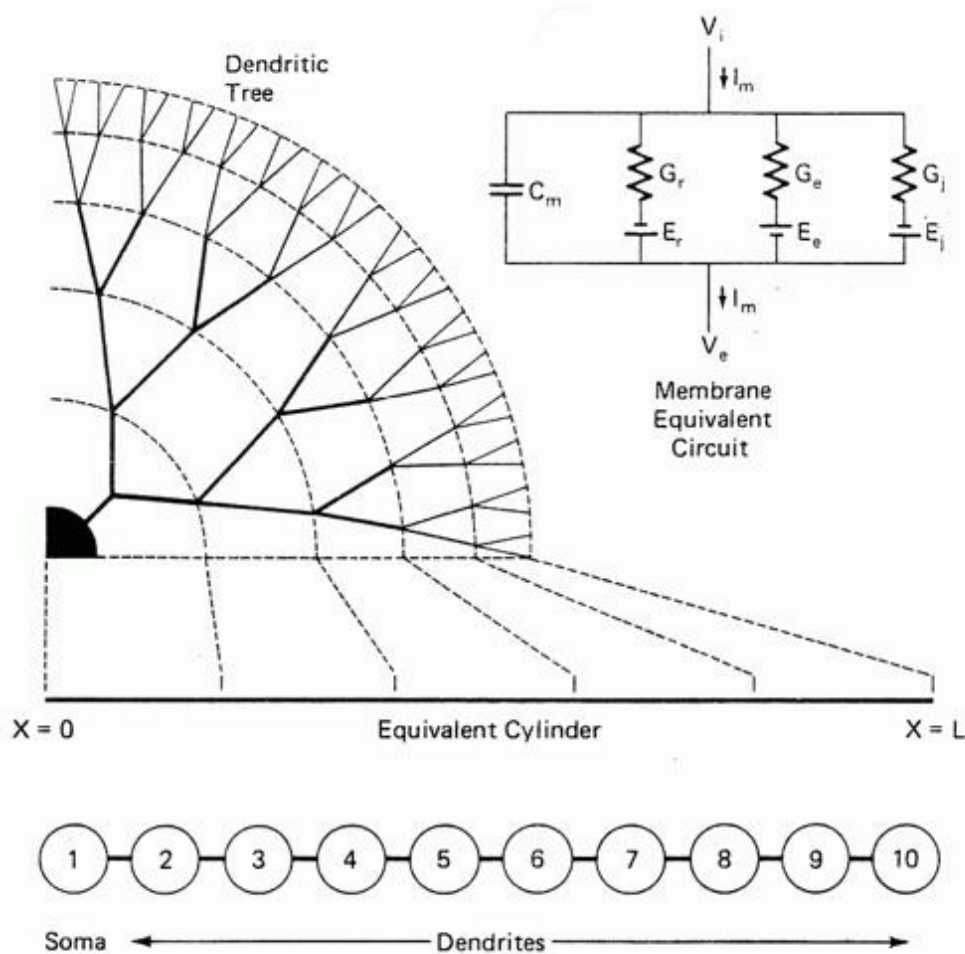


Other factors

Finite cables

Active channels

Rall model



$$R_N = \frac{\sqrt{R_m R_i / 2}}{2\pi a^{3/2}}$$

Impedance matching:

$$\text{If } a^{3/2} = d_1^{3/2} + d_2^{3/2}$$

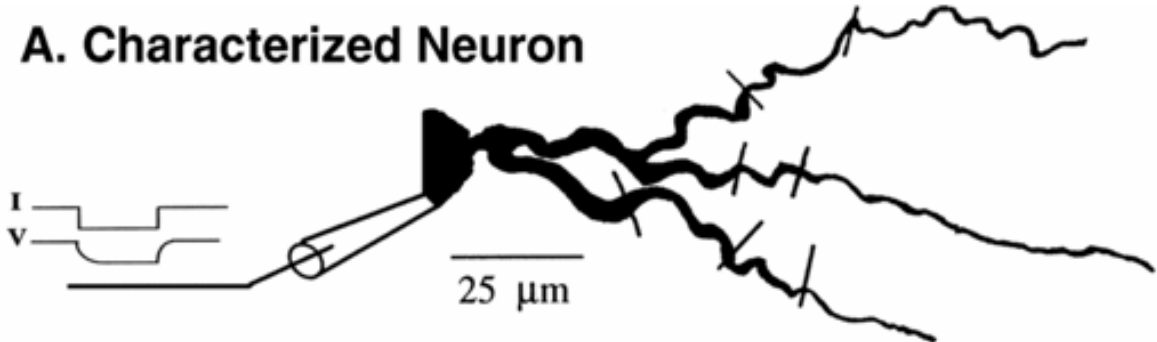
can collapse to an equivalent cylinder with length given by electrotonic length

Active cables

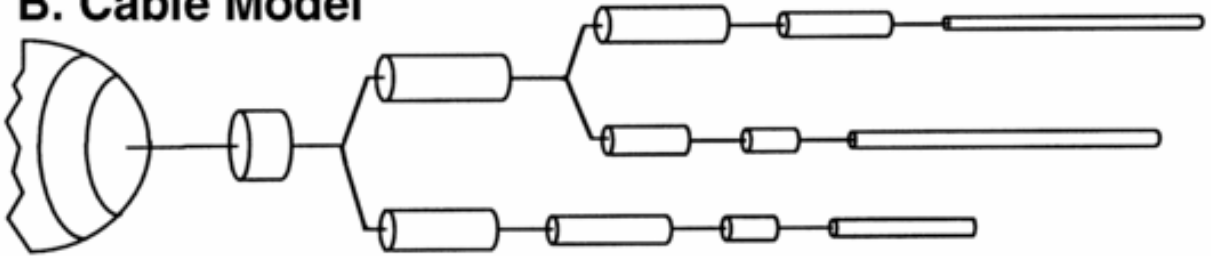
New cable equation for each dendritic compartment

Who'll be my Rall model, now that my Rall model is gone

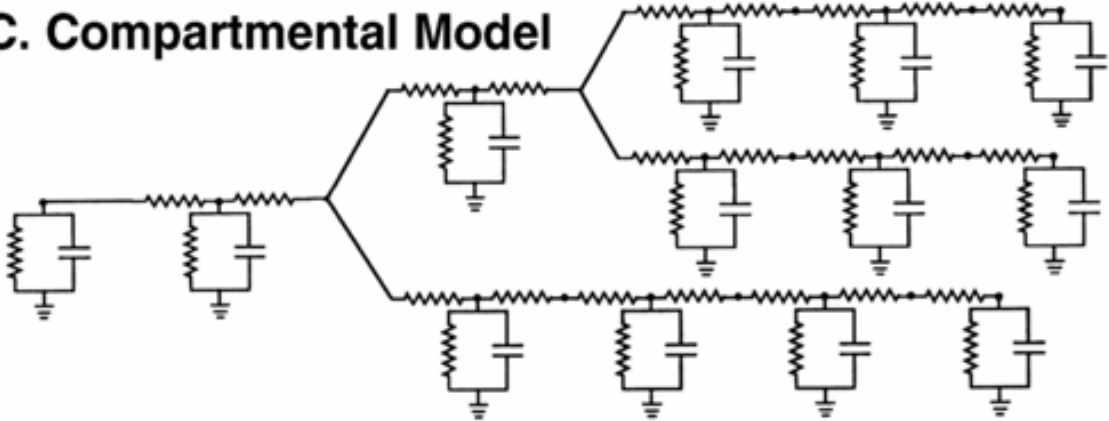
A. Characterized Neuron

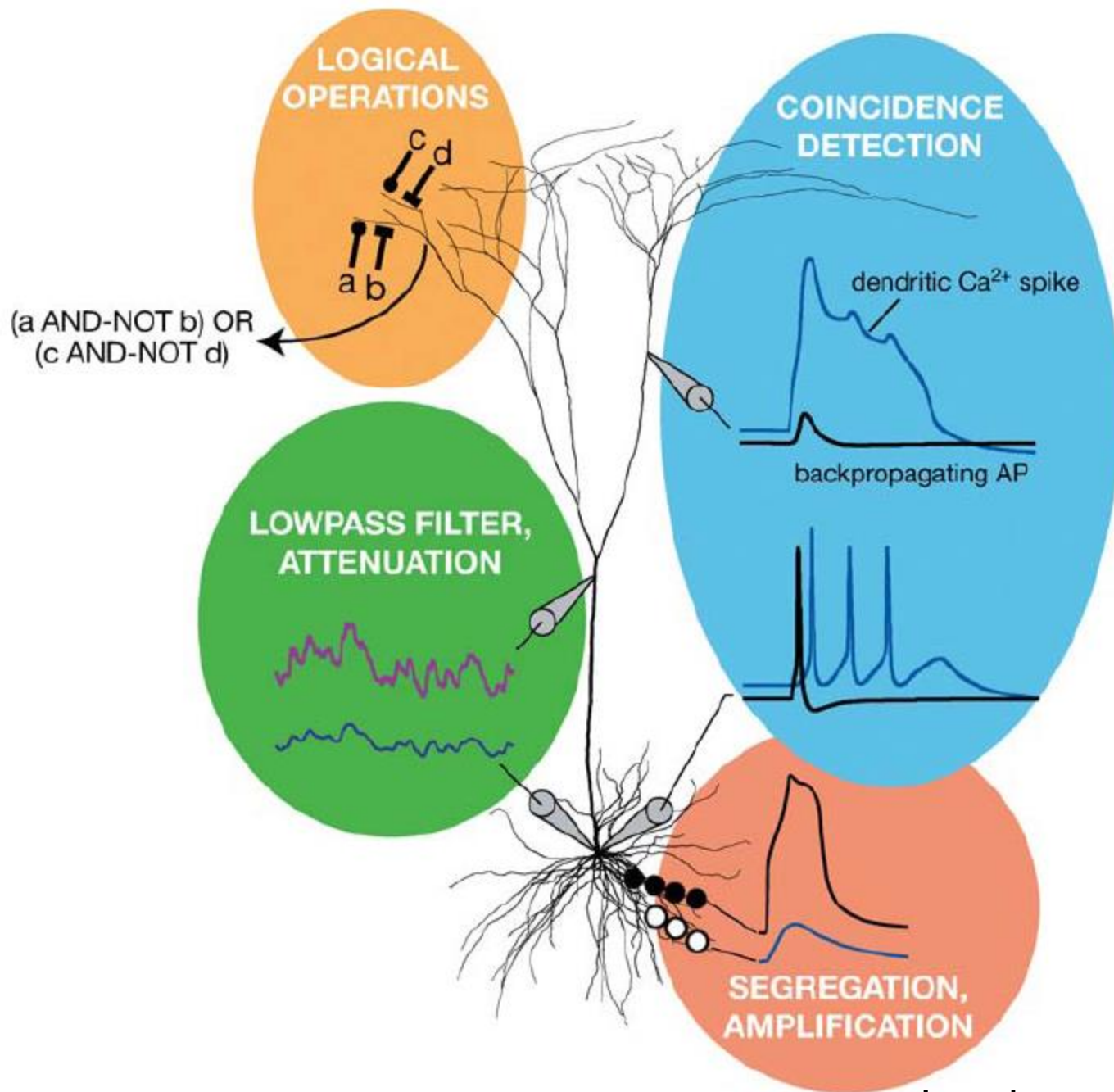


B. Cable Model



C. Compartmental Model





Enthusiastically recommended references

- **Johnson and Wu, *Foundations of Cellular Physiology*, Chap 4**
The classic textbook of biophysics and neurophysiology: lots of problems to work through. Good for HH, ion channels, cable theory.
- **Koch, *Biophysics of Computation***
Insightful compendium of ion channel contributions to neuronal computation
- **Izhikevich, *Dynamical Systems in Neuroscience***
An excellent primer on dynamical systems theory, applied to neuronal models
- **Magee, *Dendritic integration of excitatory synaptic input*,
Nature Reviews Neuroscience, 2000**
Review of interesting issues in dendritic integration
- **London and Häusser, *Dendritic Computation*,
Annual Reviews in Neuroscience, 2005**
Review of the possible computational space of dendritic processing